



David Deutsch



Richard Jozsa



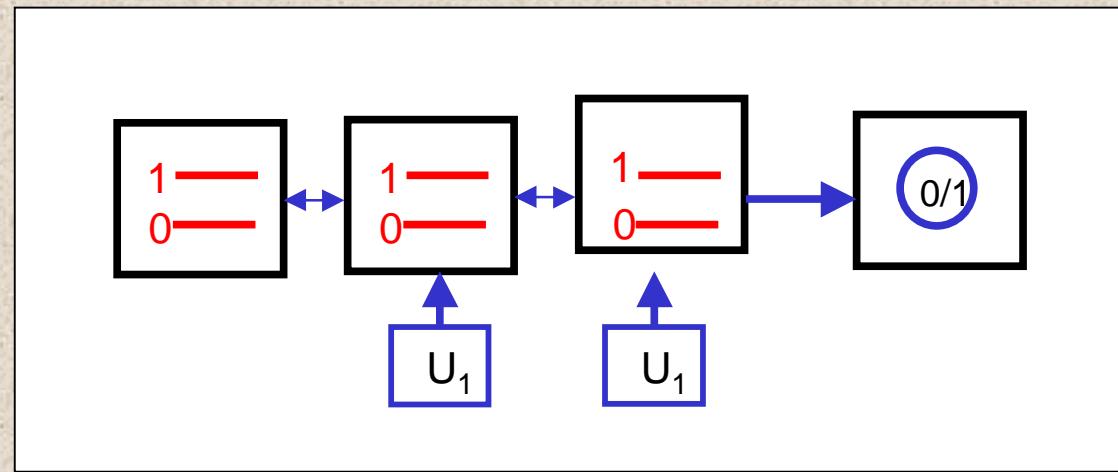
Peter Shor



Lov Grover

powerful quantum
algorithms . . .

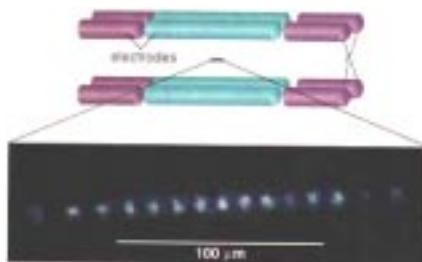
...demanding powerful
quantum hardware



micro versus macro

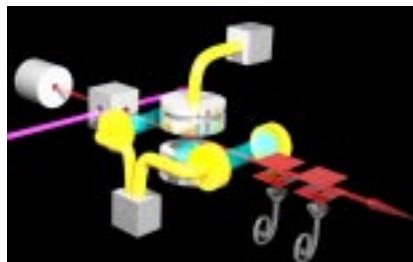
Quantum optics

Trapped ions



(NIST, Innsbrück...)

Atoms in cavity



(ENS Paris)

Quantum,
but not easily scalable

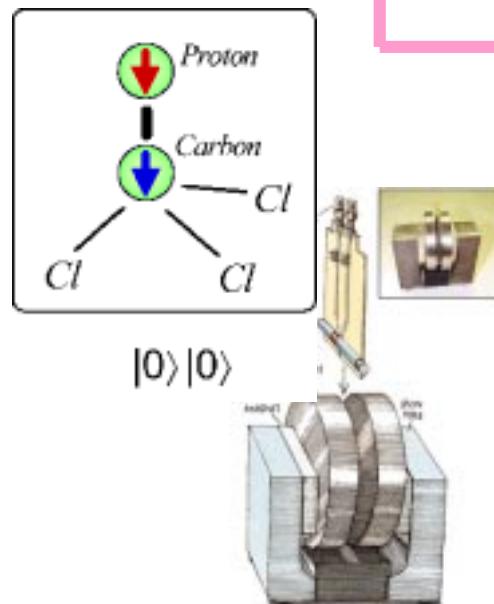
NMR

e / L He
atoms

on
chip

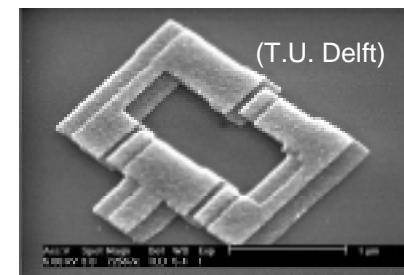
Solid-state devices

semiconductors
P in Si
spintronics →



(Oxford, Stanford,
IBM,MIT...)

superconducting circuits



SQUBIT 1-2

Chalmers U.
TU Delft
PTB
Quantronics

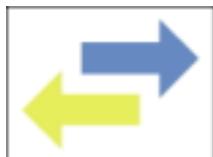
CRTBT
NEC
NTT
NIST Boulder
Stony Brook
Kansas U.
Yale
UCB
...

Scalable,
but not easily quantum

Operation of a solid-state quantum-bit

**QUANTUM
ELECTRONICS GROUP**

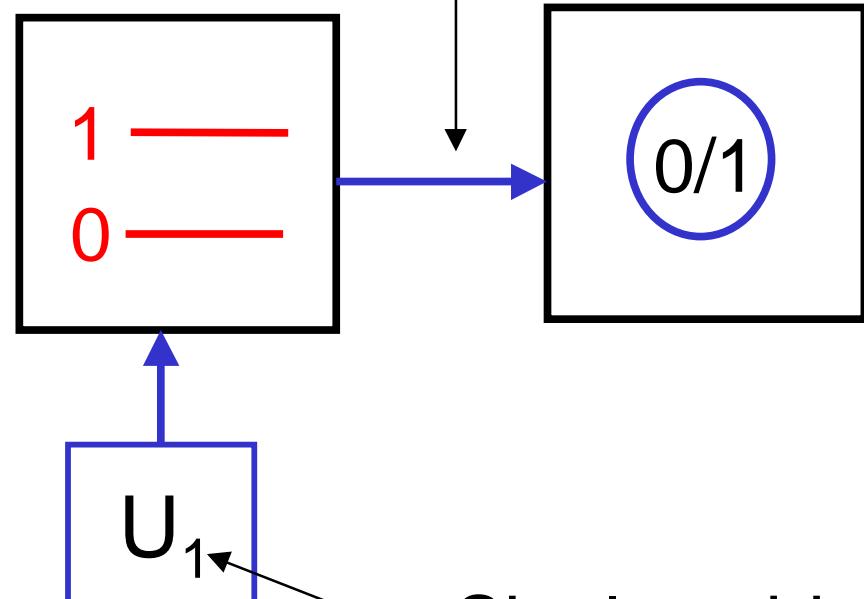
CEA-Saclay, FRANCE



**SQUBIT
collaboration**

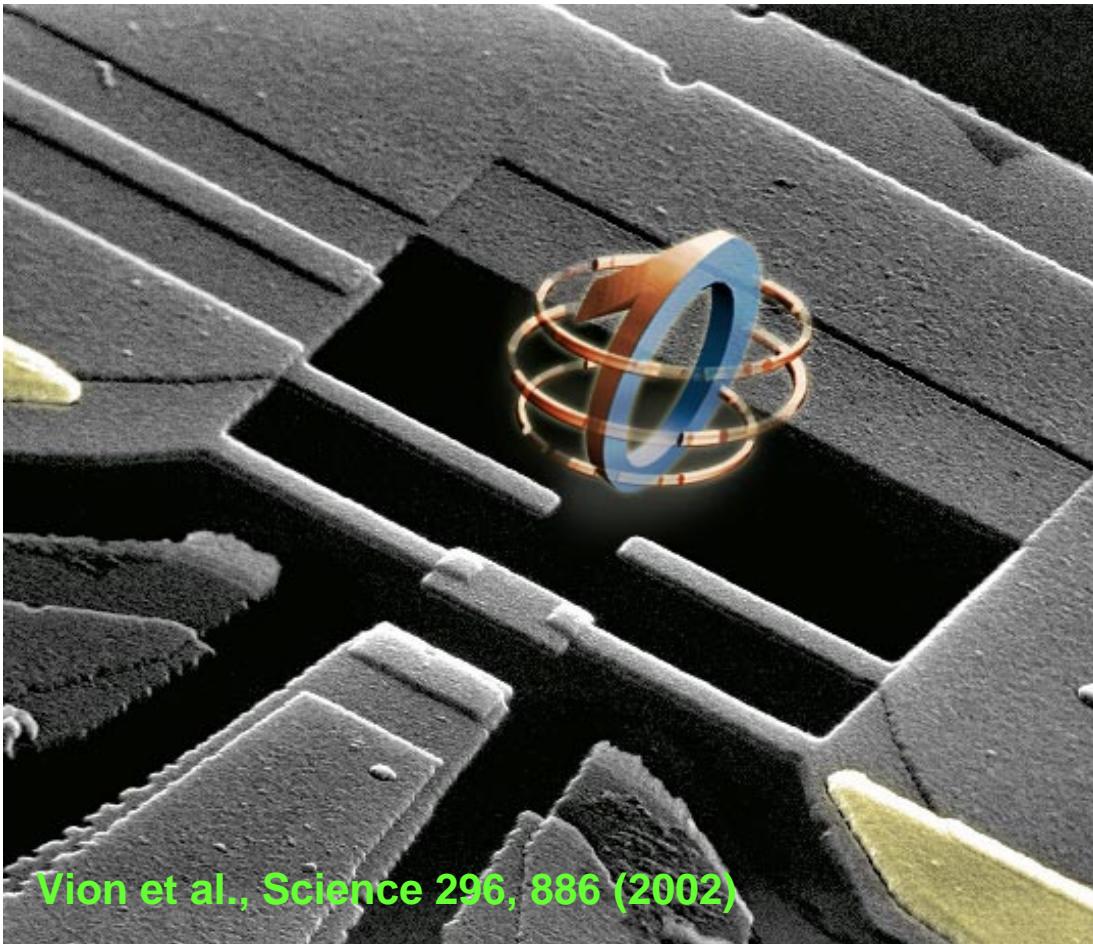


readout



**Single qubit
operations**

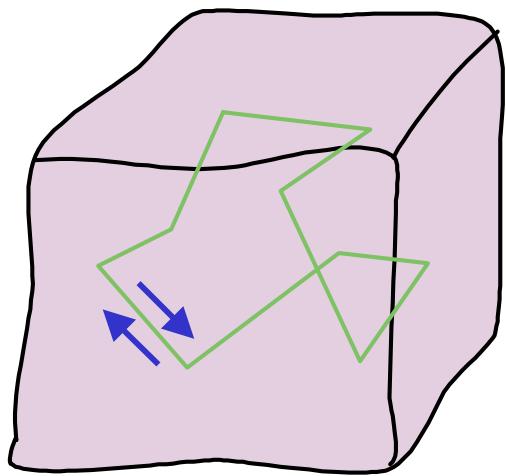
Operation of a solid-state quantum bit



Vion et al., Science 296, 886 (2002)

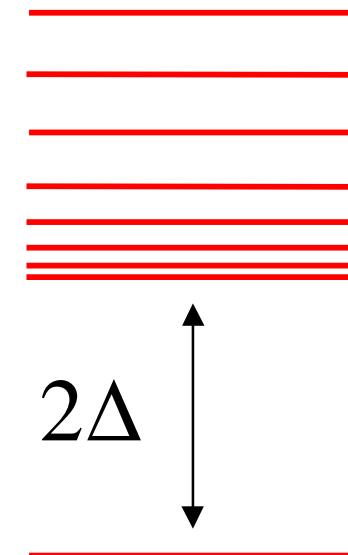
Quantronium superconducting circuit

Why superconductivity ?



All states paired

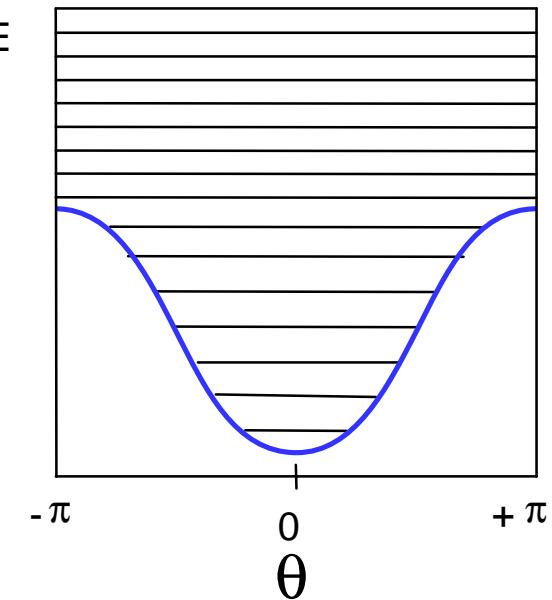
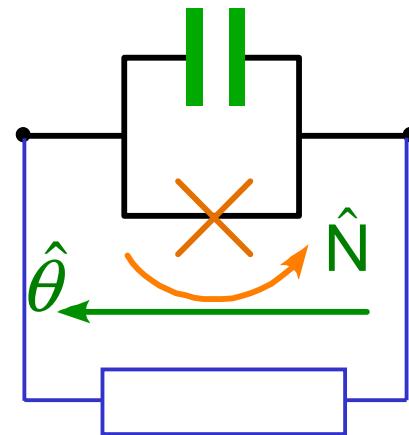
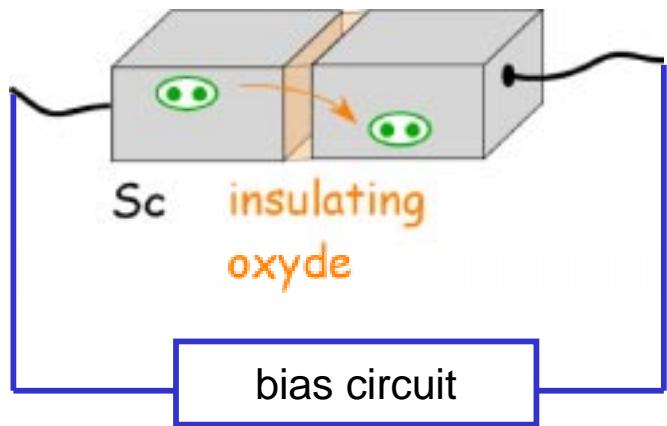
$$|\langle \circ \circ \rangle + \langle \bullet \bullet \rangle\rangle$$



Superconducting Condensate
Ground state

The Josephson junction

A single degree of freedom $[\hat{N}, \hat{\theta}] = i$



Josephson hamiltonian:

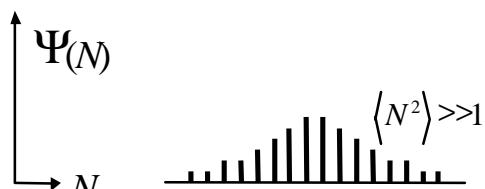
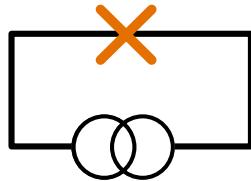
$$H_J = -E_J \cos \hat{\theta} = -\frac{E_J}{2} \sum_N |N\rangle \langle N+1| + |N+1\rangle \langle N|$$

full hamiltonian:

$$H = H_J + H_{elm}$$

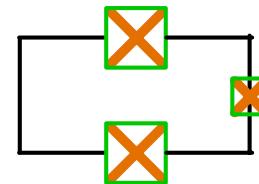
Josephson qubits

Current-biased
large junction



Phase state

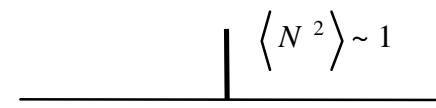
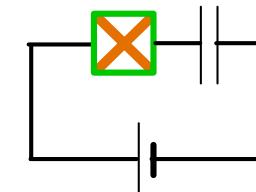
Medium-size
junctions in a loop



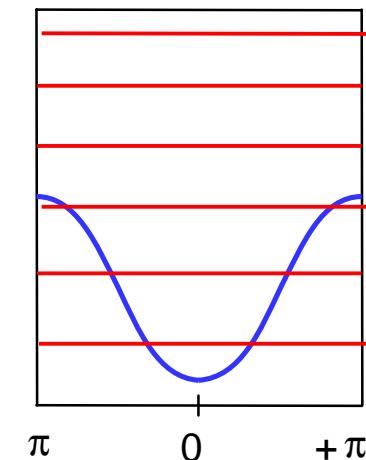
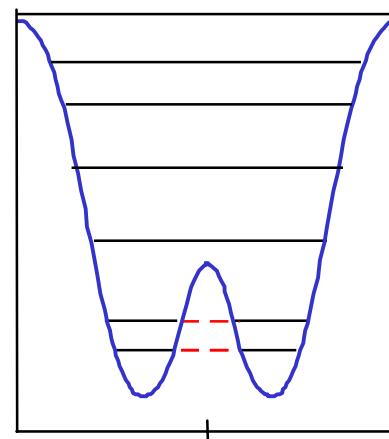
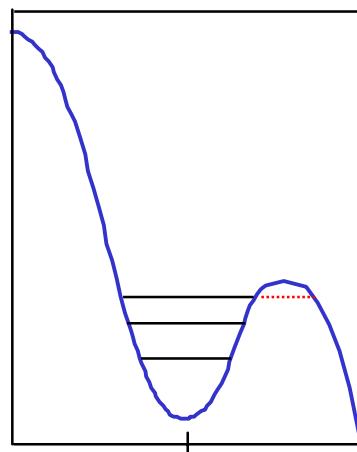
phase

charge

Small junction
in a box geometry

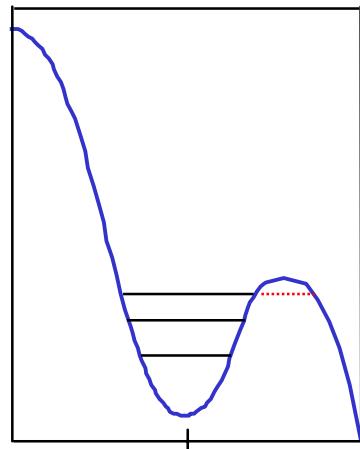
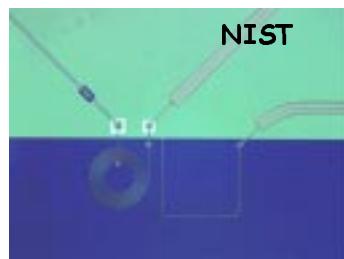
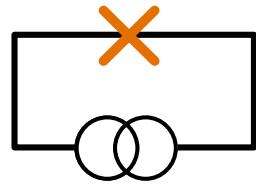


charge state

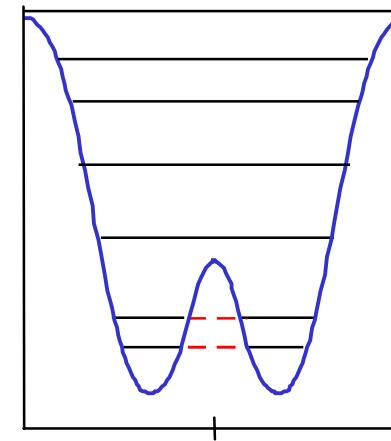
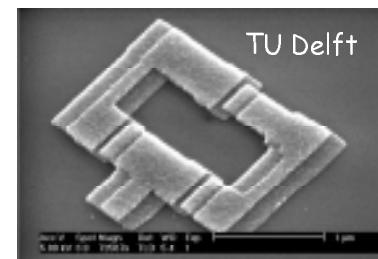
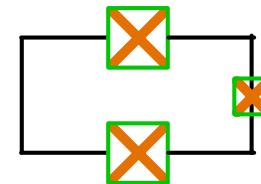


Josephson qubits

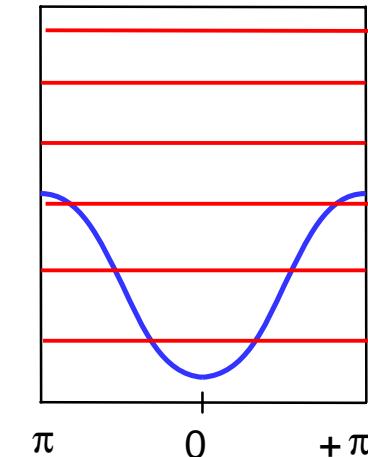
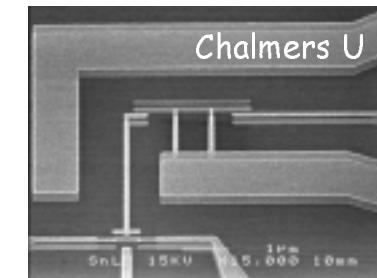
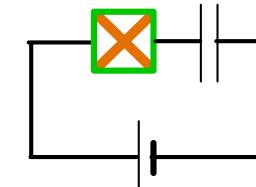
Current-biased
large junction



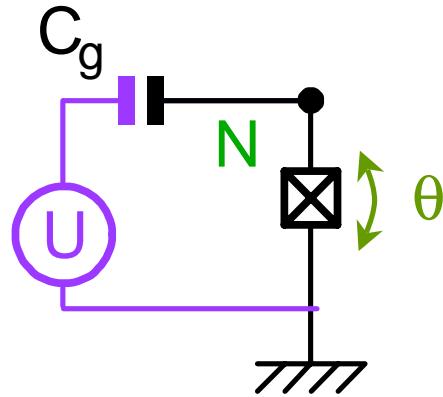
Medium-size
junctions in a loop



Small junction
in a box geometry



The Cooper pair box



1 degree of freedom

$$[\hat{N}, \hat{\theta}] = i$$

2 characteristic energies

$$E_c = \frac{(2e)^2}{2C_{\text{island}}} \quad E_J = \frac{\Delta h}{8 e^2 R_T}$$

"potential"

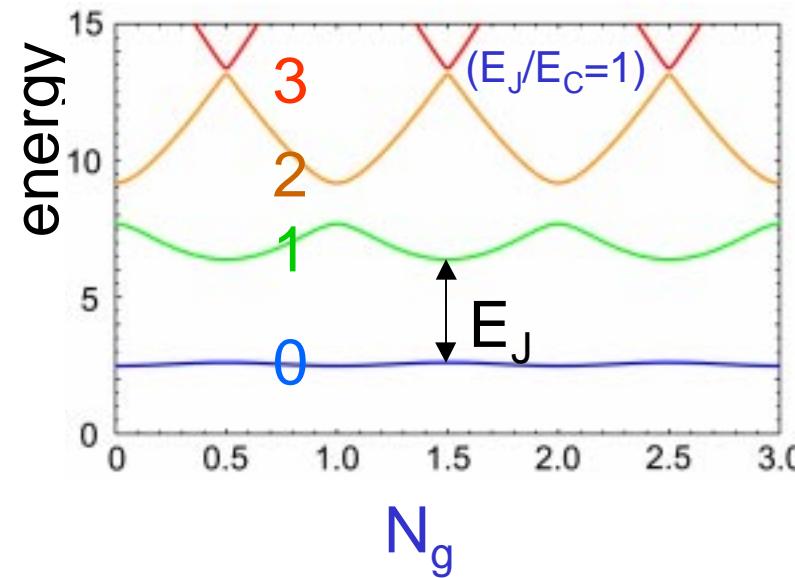
"kinetic"

Hamiltonian

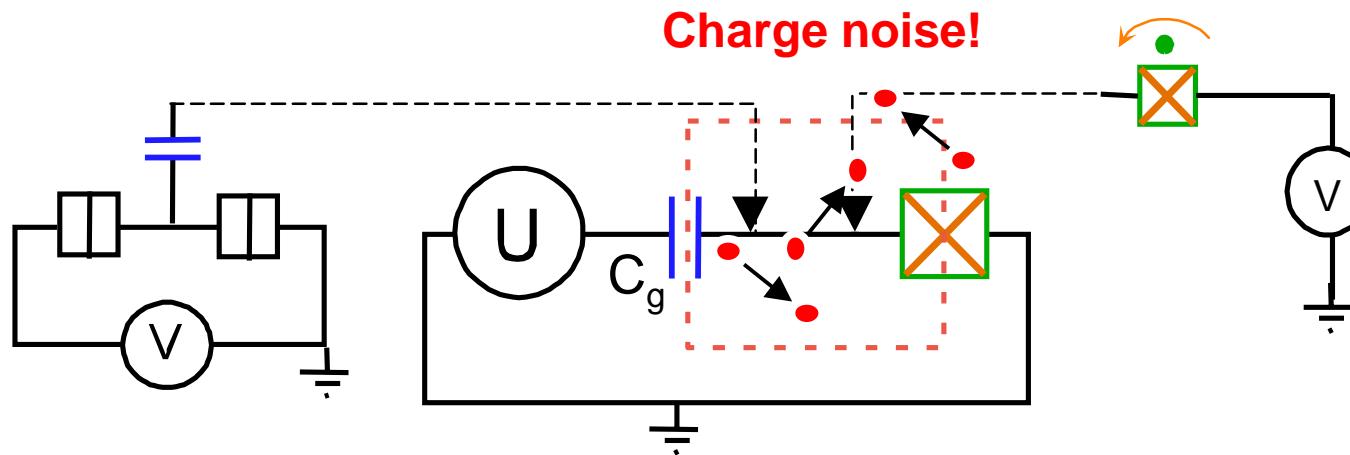
$$\hat{H} = E_c (\hat{N} - N_g)^2 - E_J \cos \hat{\theta}$$

1 control knob $C_g U / 2e$

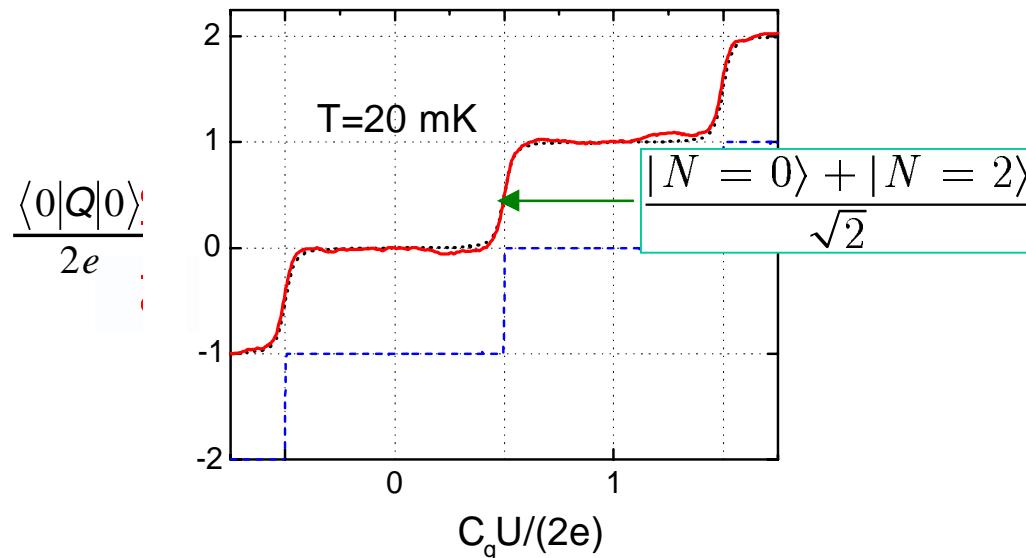
Spectrum



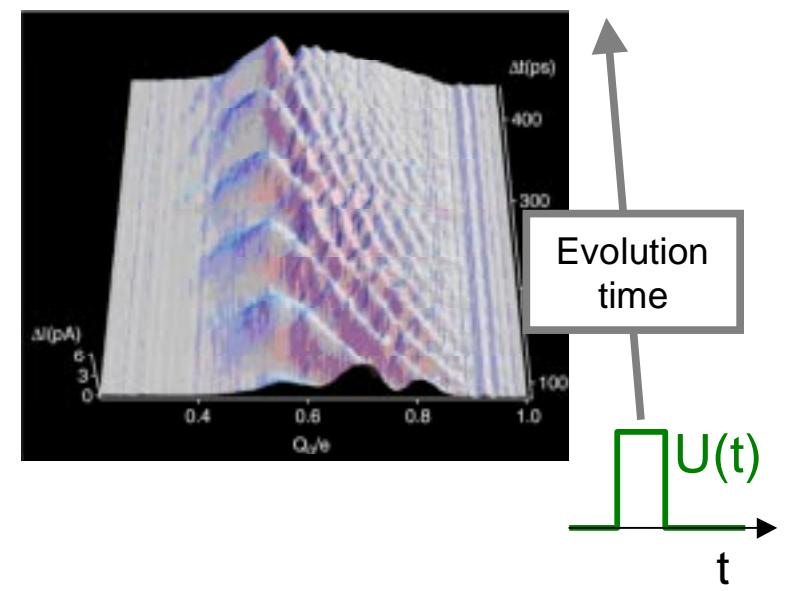
Measuring the Cooper pair box



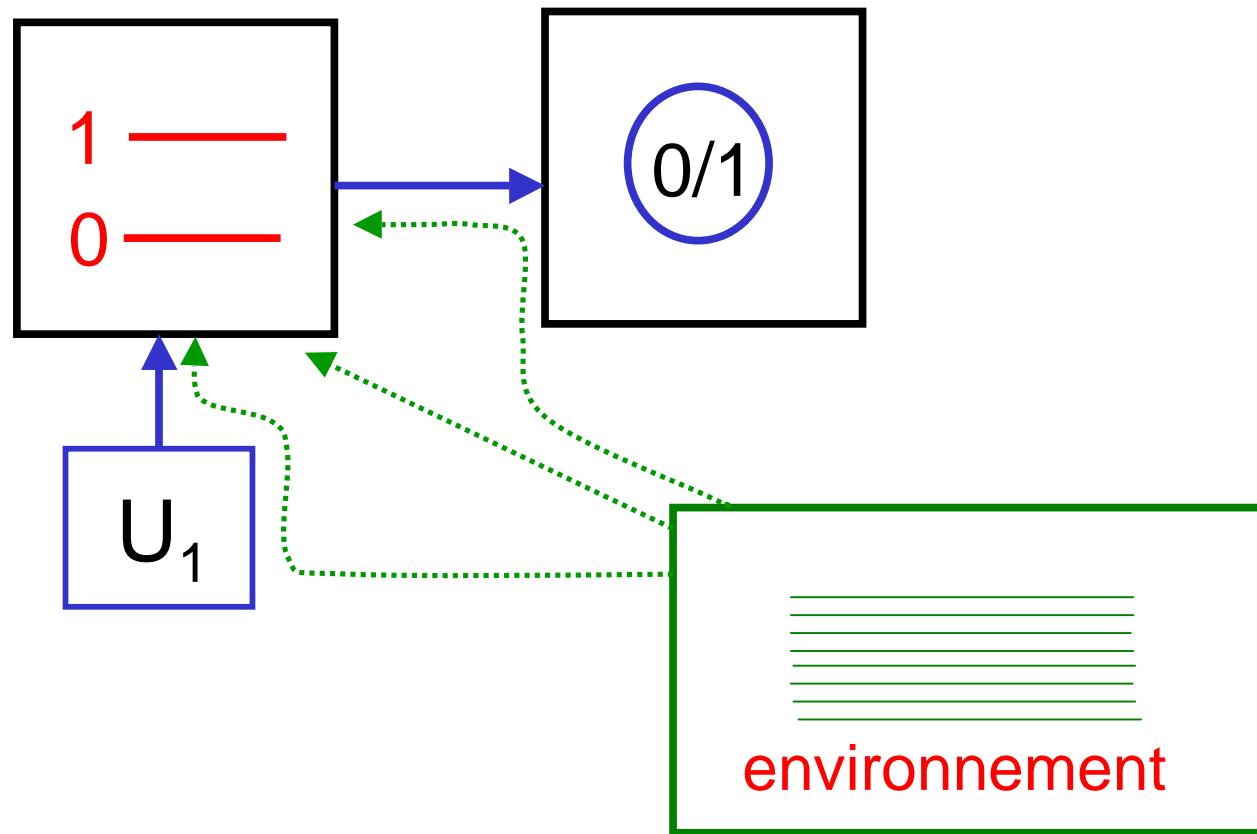
1996 charge of ground state $|0\rangle$
 (Bouchiat et al., Quantronics)



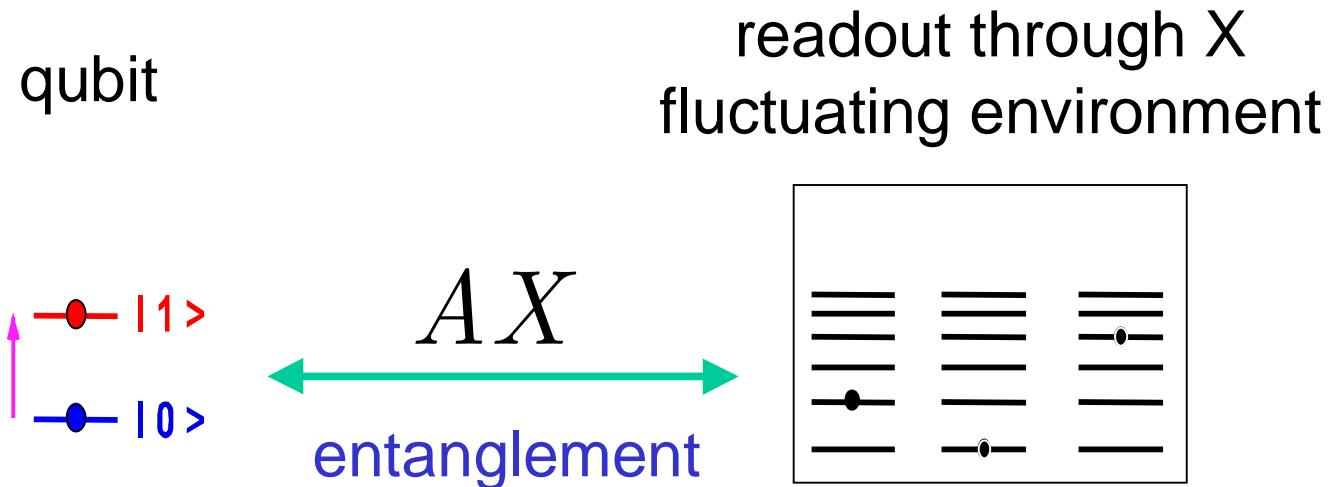
1999 coherent superpositions $\alpha|0\rangle + \beta|1\rangle$
 (Nakamura, Pashkin & Tsai, NEC)



decoherence and readout



decoherence and readout



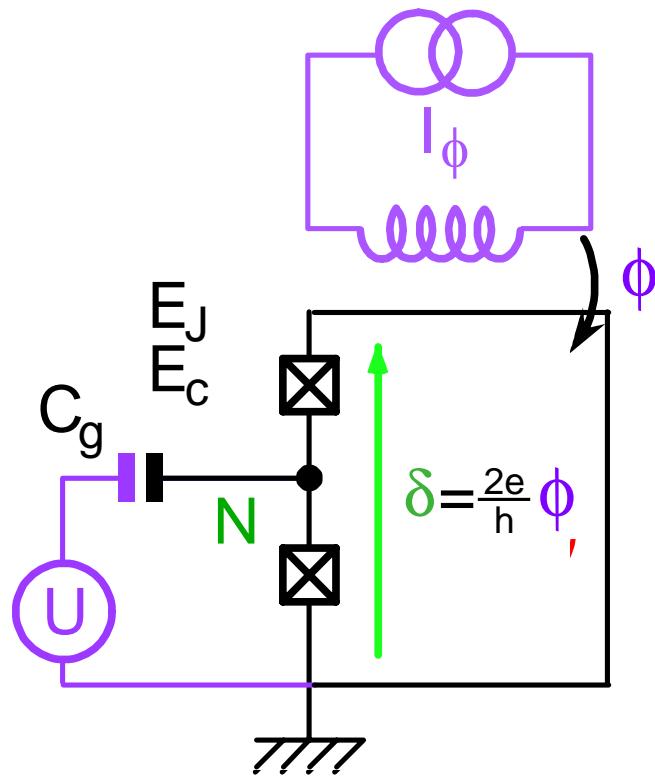
Signal : $[\langle 1|A|1 \rangle - \langle 0|A|0 \rangle] = h \frac{\partial \nu_{01}}{\partial X}$

Dephasing : $\delta X(t) \longrightarrow \delta \nu_{01}(t) = \frac{\partial \nu_{01}}{\partial X} \delta X(t)$

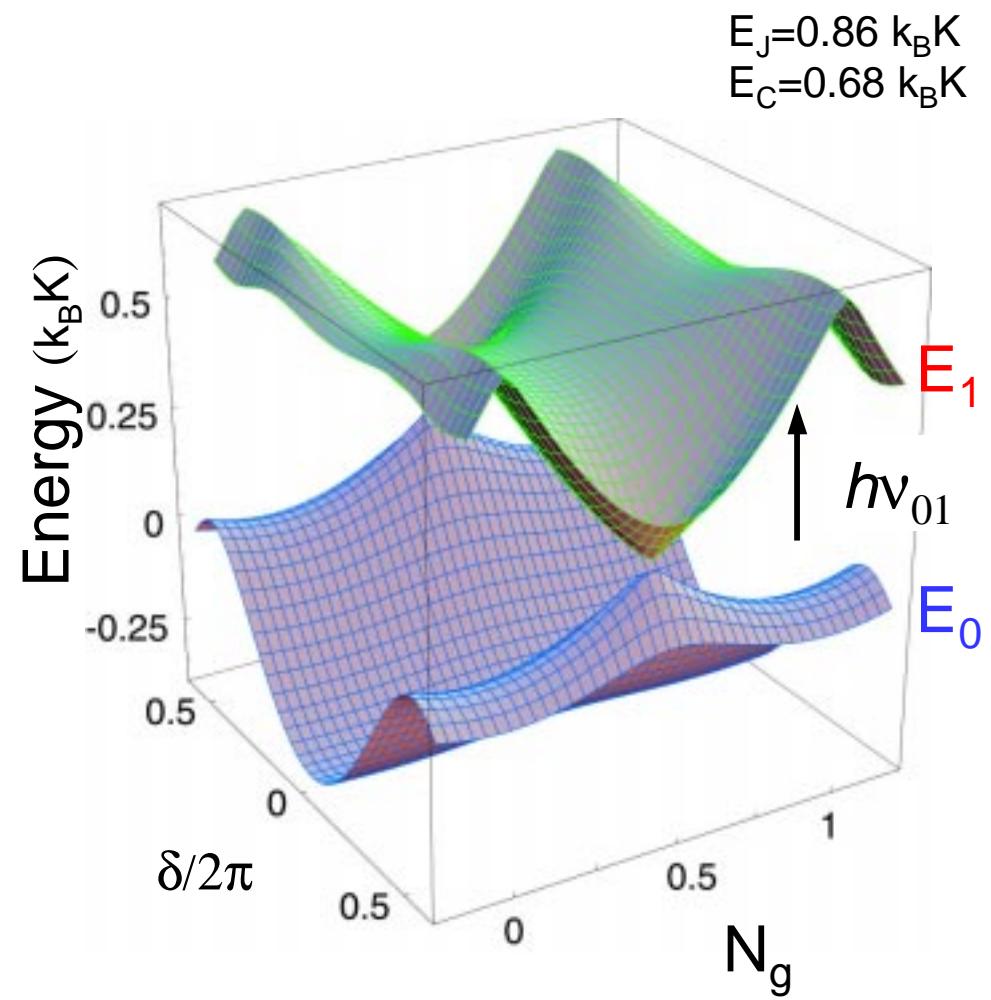
Readout + environment

Move adiabatically **then** readout

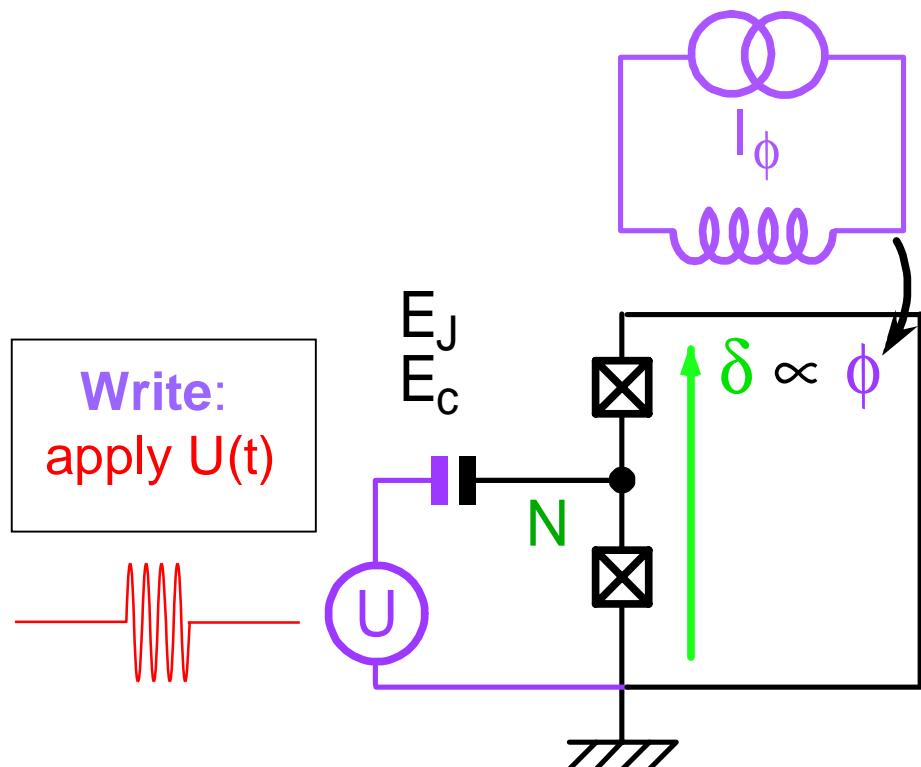
The Quantronium: a split junction Cooper pair box



$$E_J^{\text{eff}}(\delta) = E_J \cos \frac{\delta}{2}$$



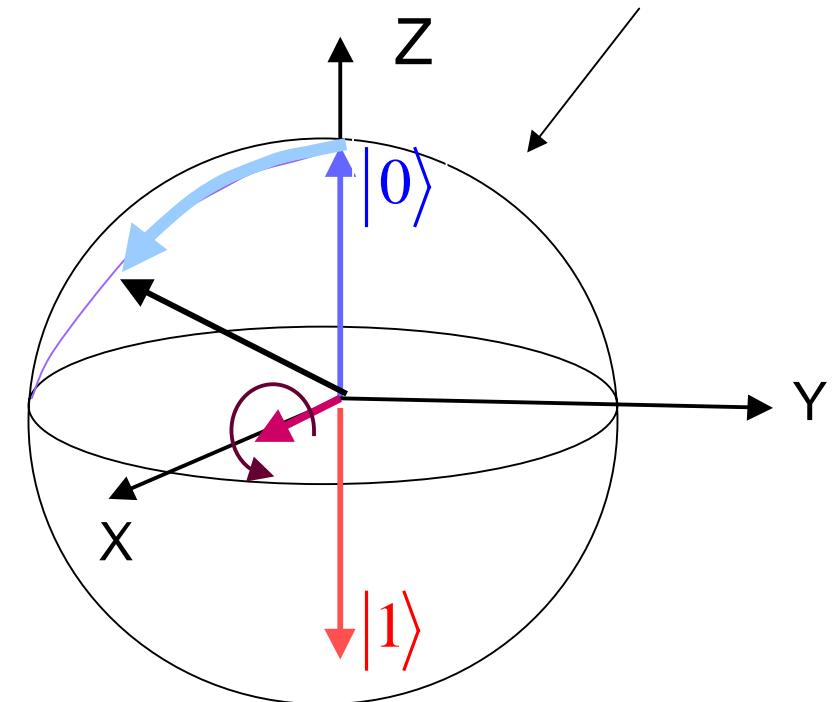
State manipulation using the charge port



Microwave drive at

$$\nu_{\mu w} \approx \nu_{01}$$

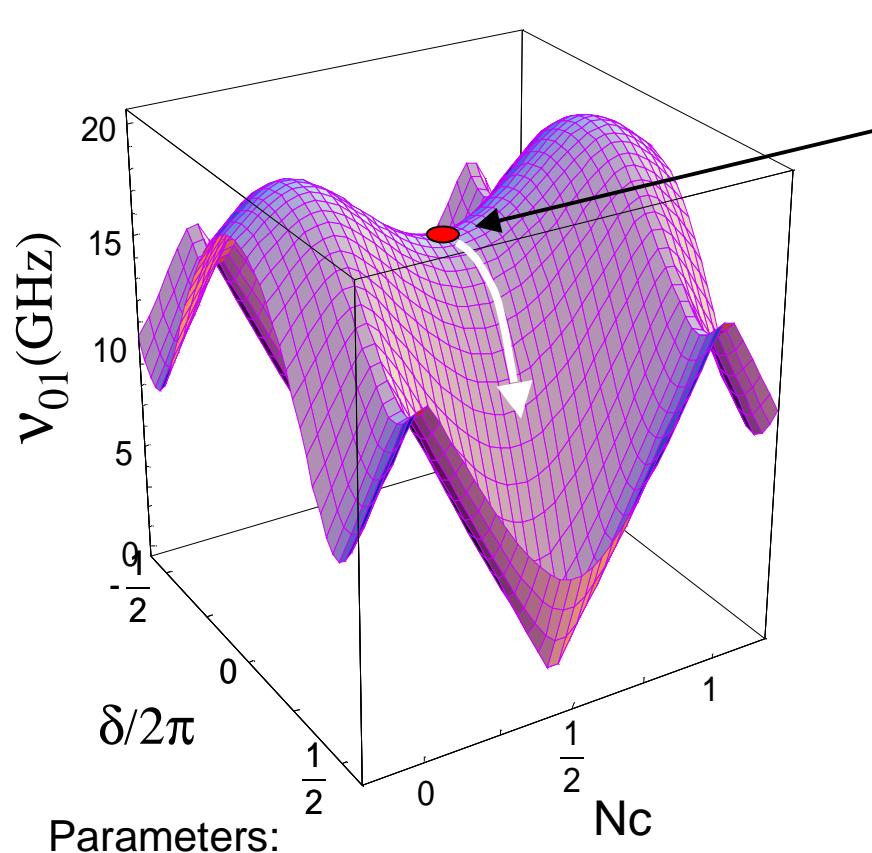
Bloch sphere representation in the rotating frame



Rabi precession

$$\omega_{\text{Rabi}} = \alpha U_{\text{RF}}$$

Decoherence and readout



Parameters:

$$E_J = 0.86$$

$$k_B K$$

$$E_C = 0.68$$

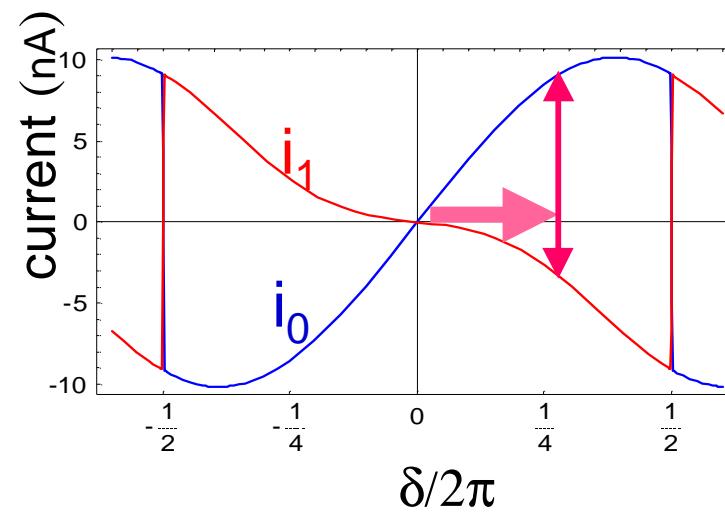
$$k_B K$$

At saddle point :

$$\frac{\partial \nu_{01}}{\partial N_g} = 0$$

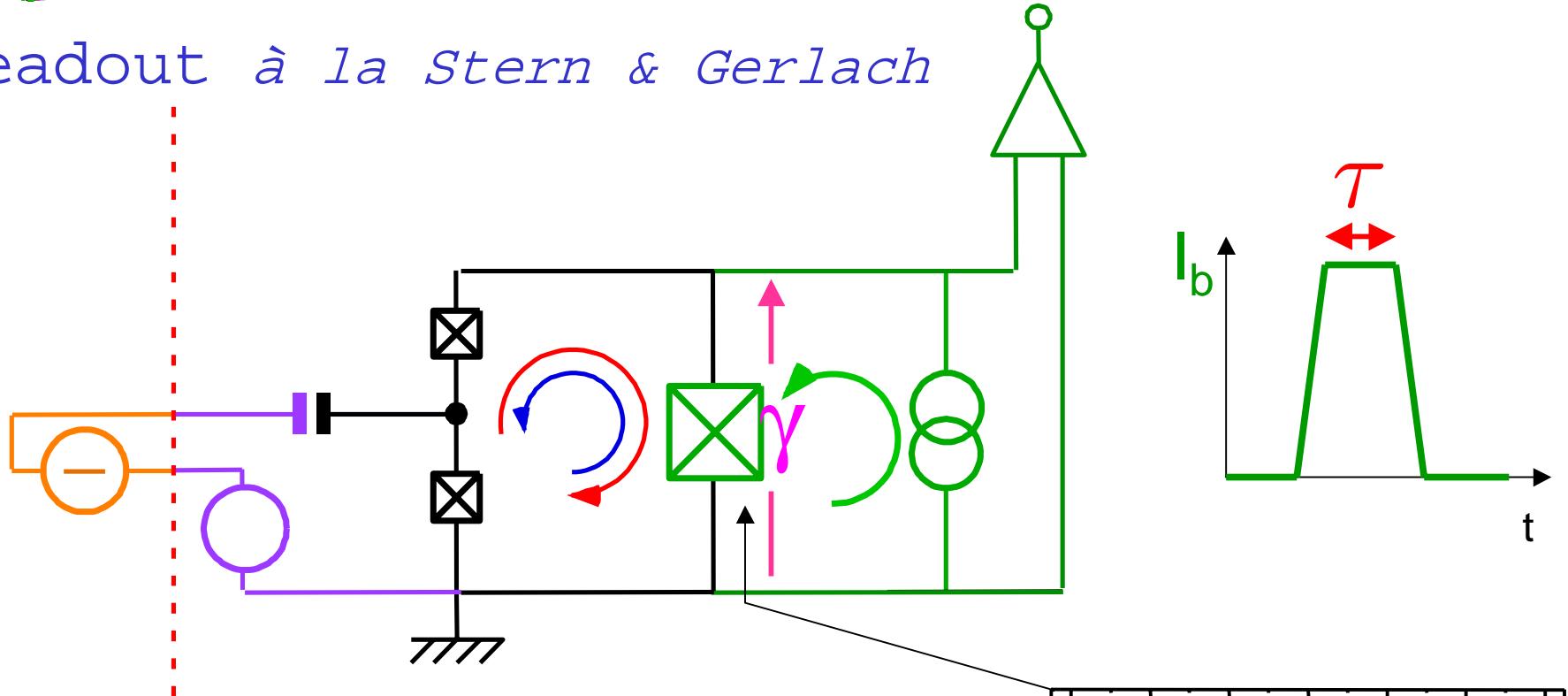
$$\frac{\partial \nu_{01}}{\partial \delta} = 0$$

no dephasing
...but no signal:



But how ?

Readout à la Stern & Gerlach

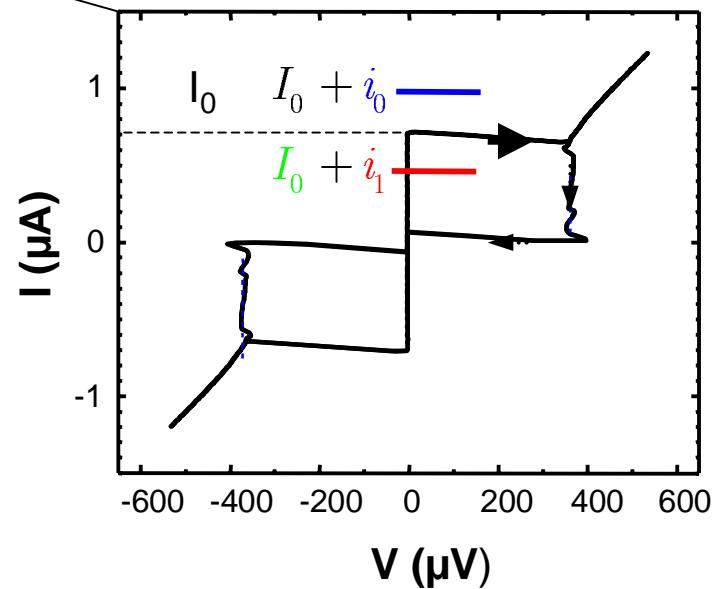


I) phase - bias :

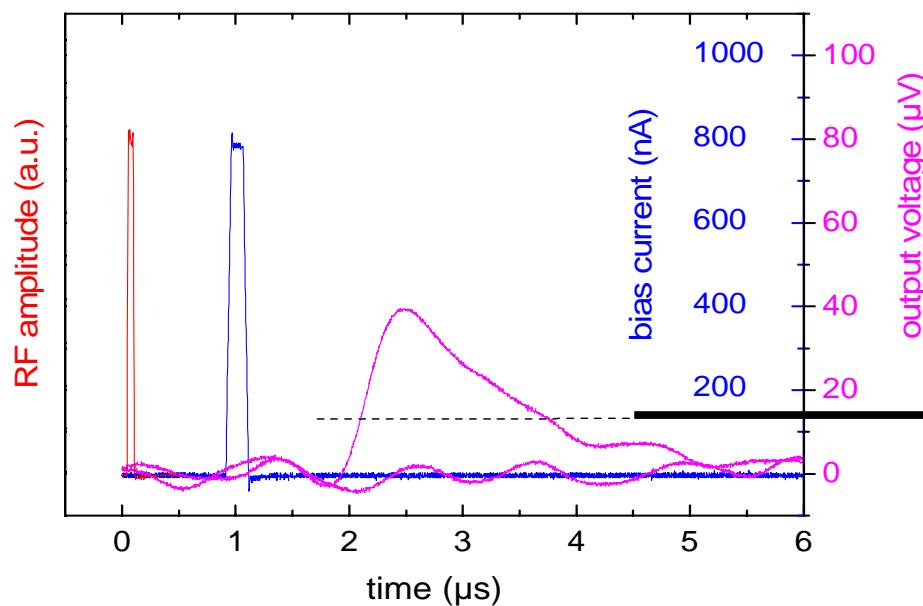
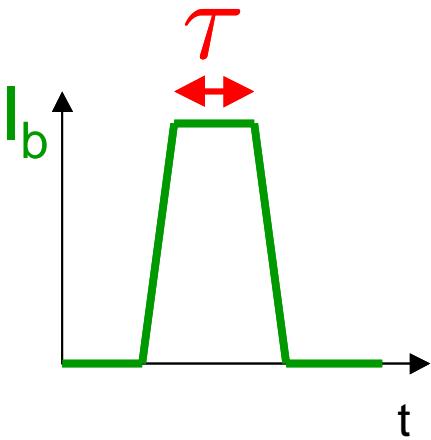
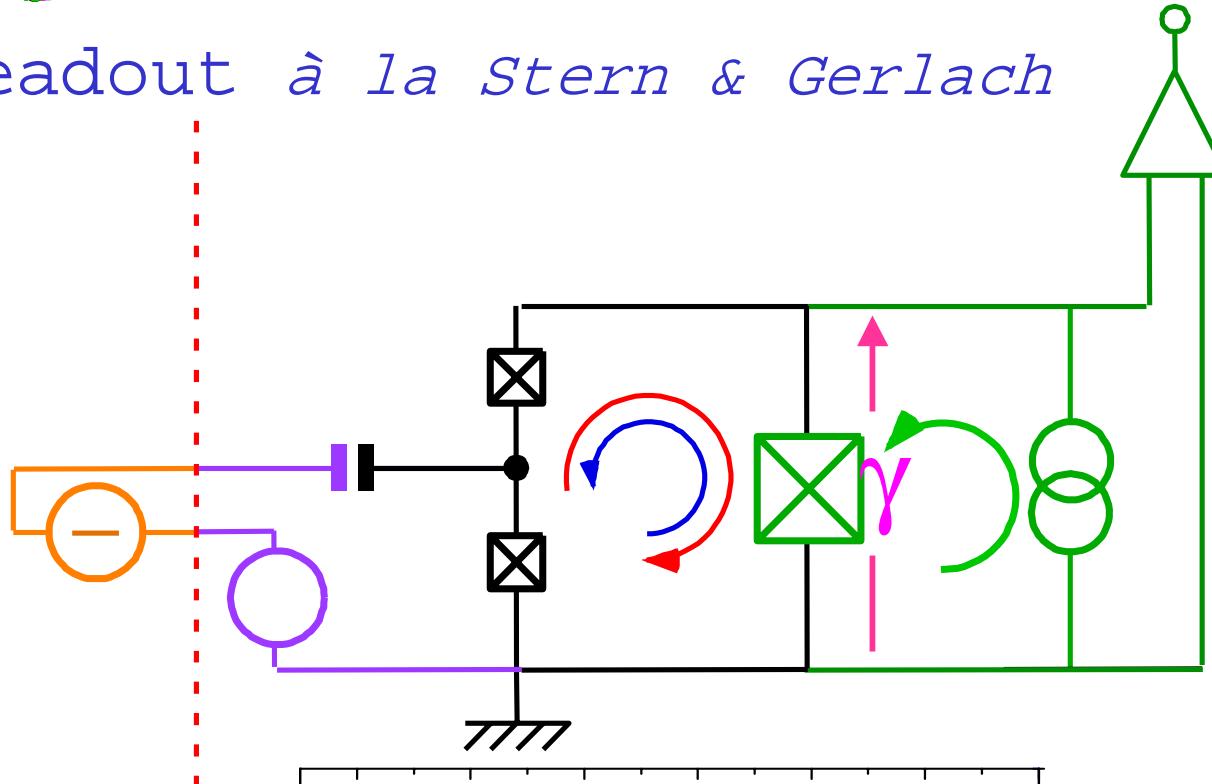
$$I_b < I_0 \quad I_b = I_0 \sin \gamma$$

II) discriminates when : $\gamma \approx \pi / 2$

$$I_b \approx I_0 \quad I_b + i_1 < I_0 < I_b + i_0$$



Readout à la Stern & Gerlach

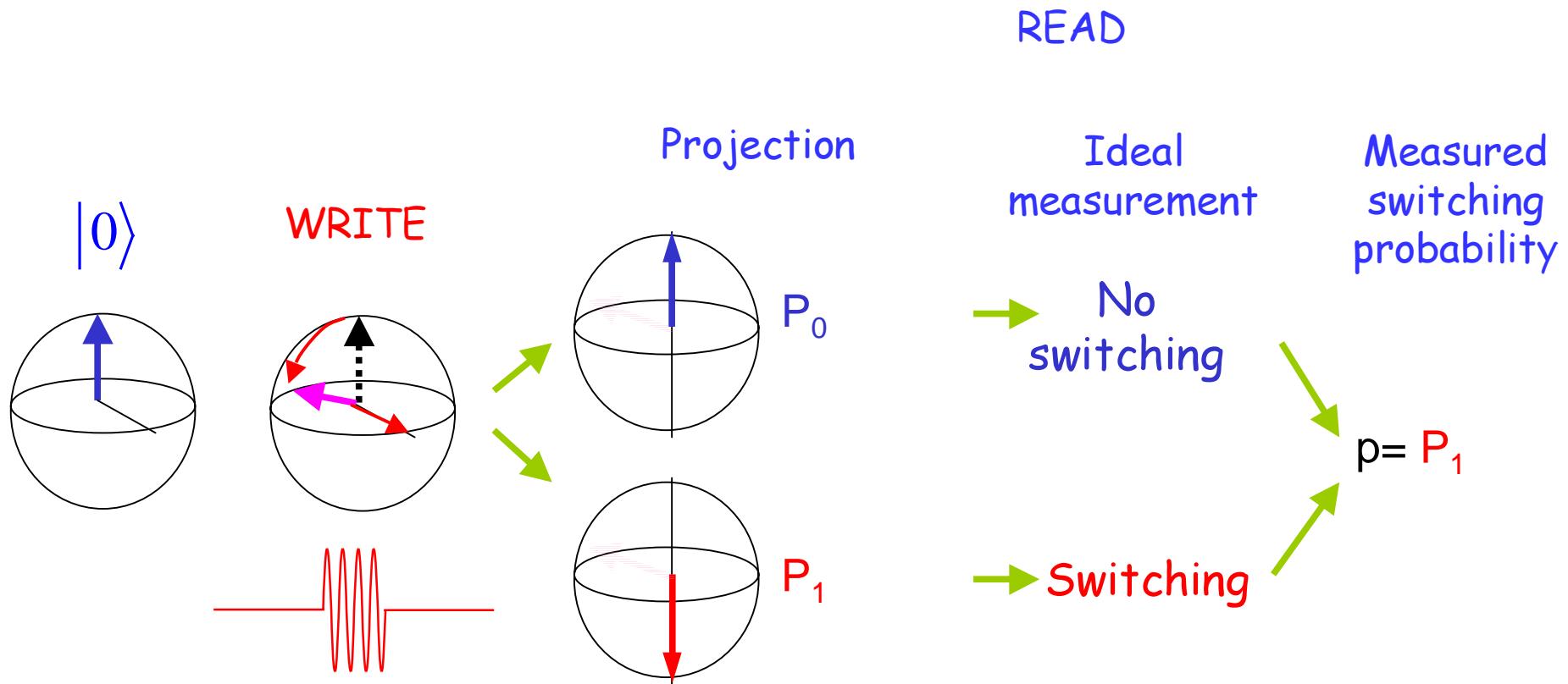


1

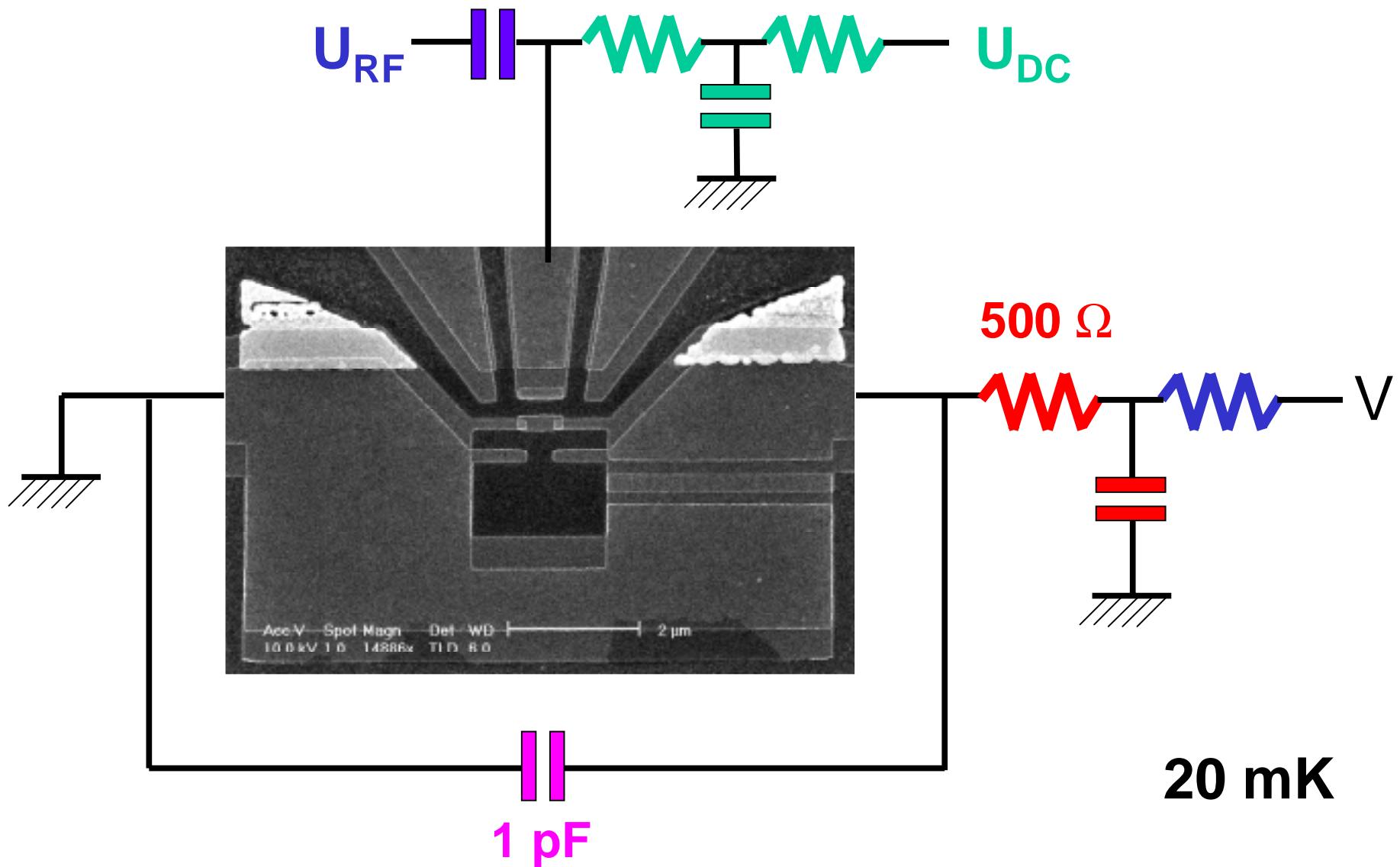
0

$\Gamma_1 \tau \gg 1$
 $\Gamma_0 \tau \ll 1$

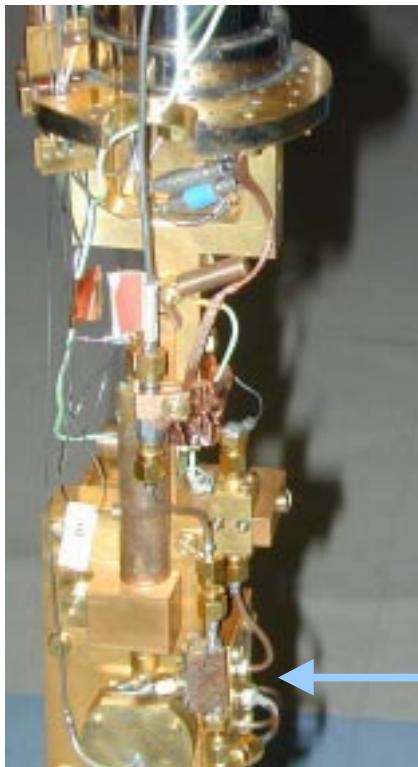
Preparation and ideal readout



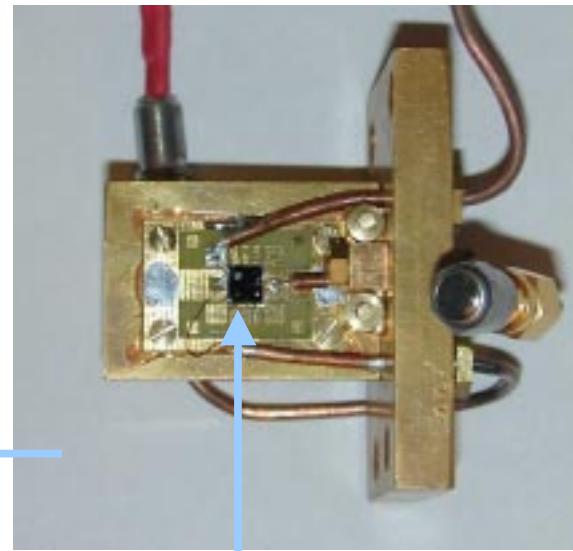
Implementation



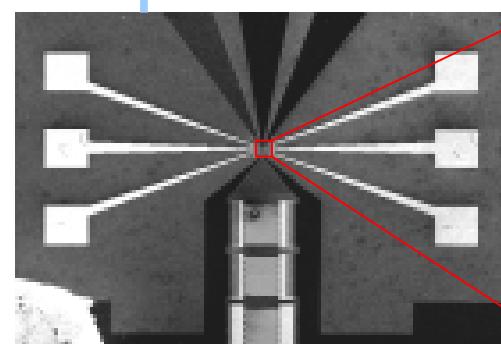
Experimental set-up



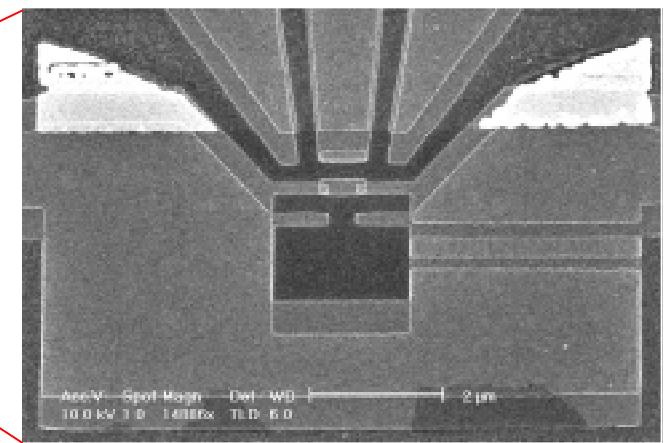
Dilution fridge
20 mK



p.c.b.

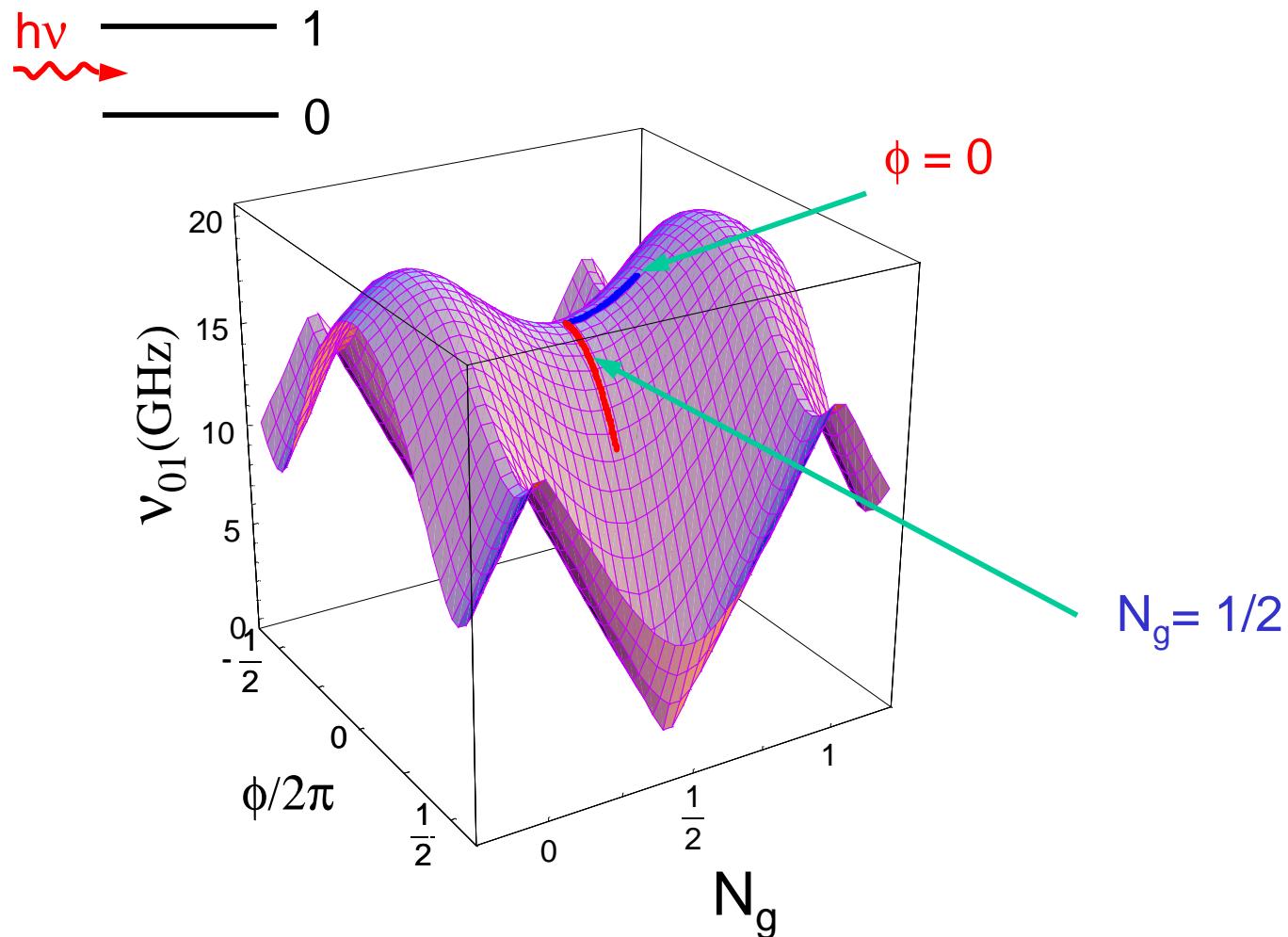


chip

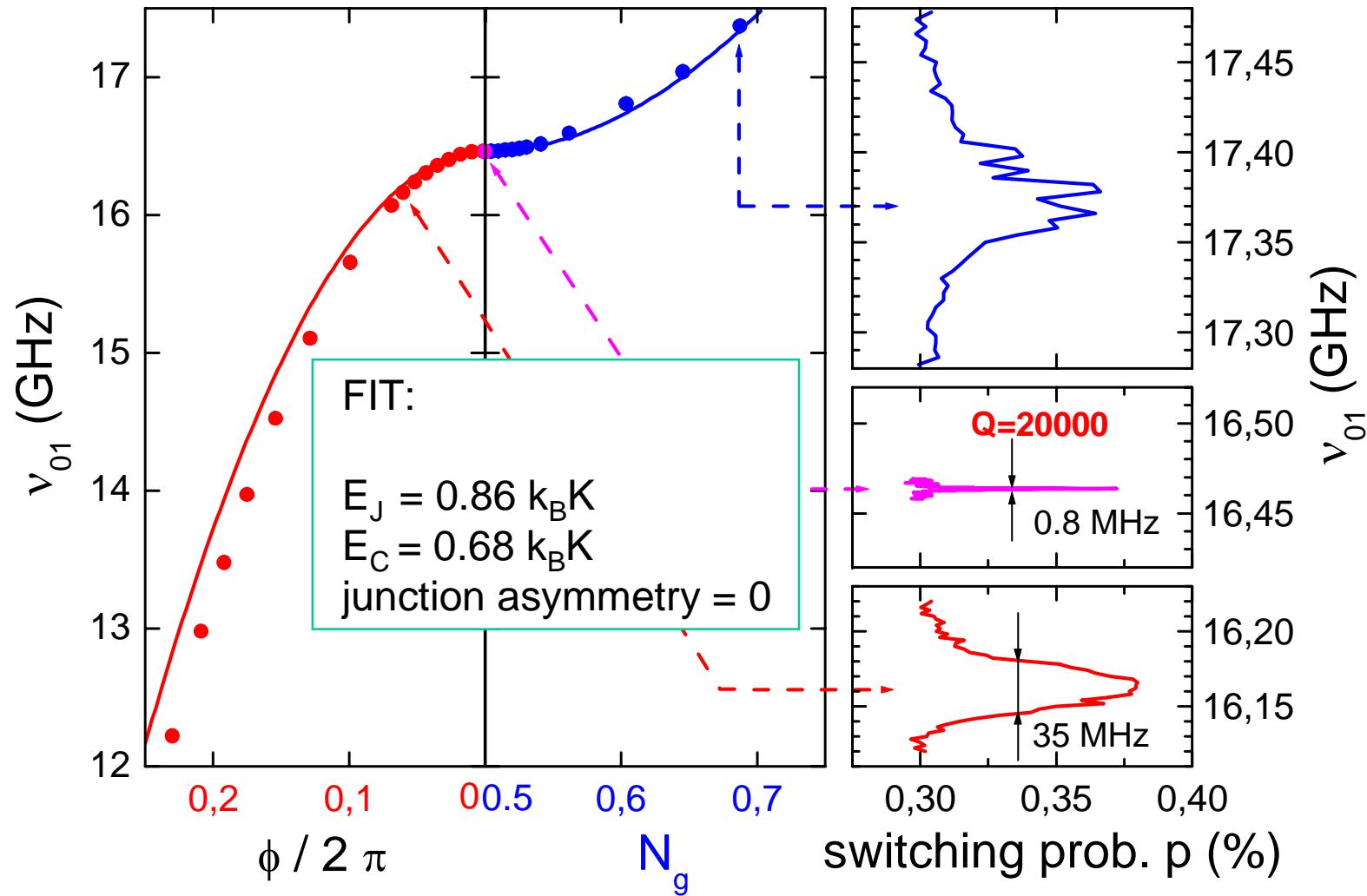


Level spectroscopy

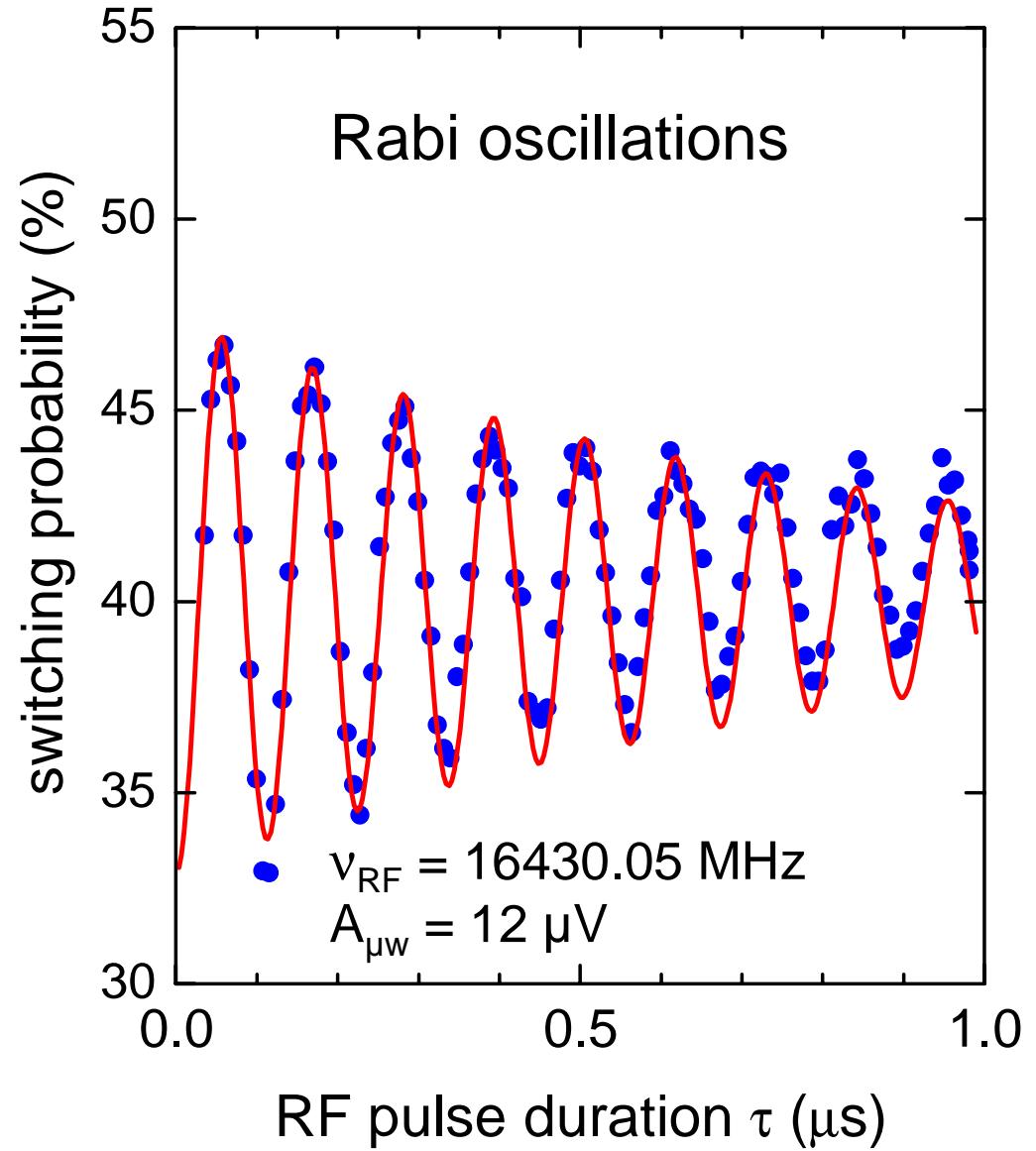
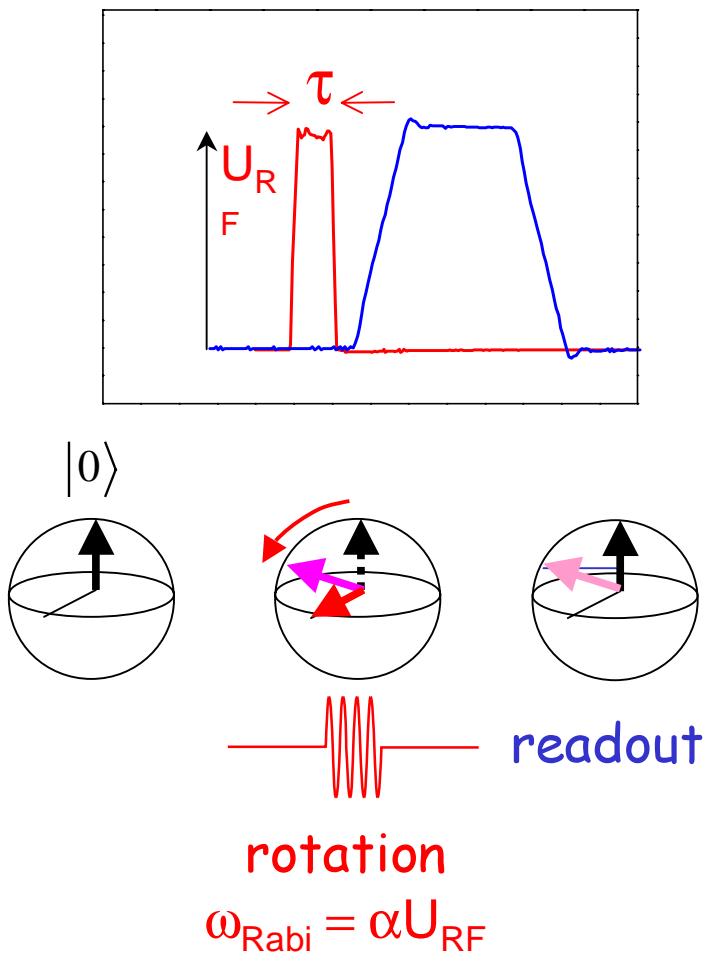
Level spectroscopy $\nu_{01}(N_g, \phi/2\pi)$



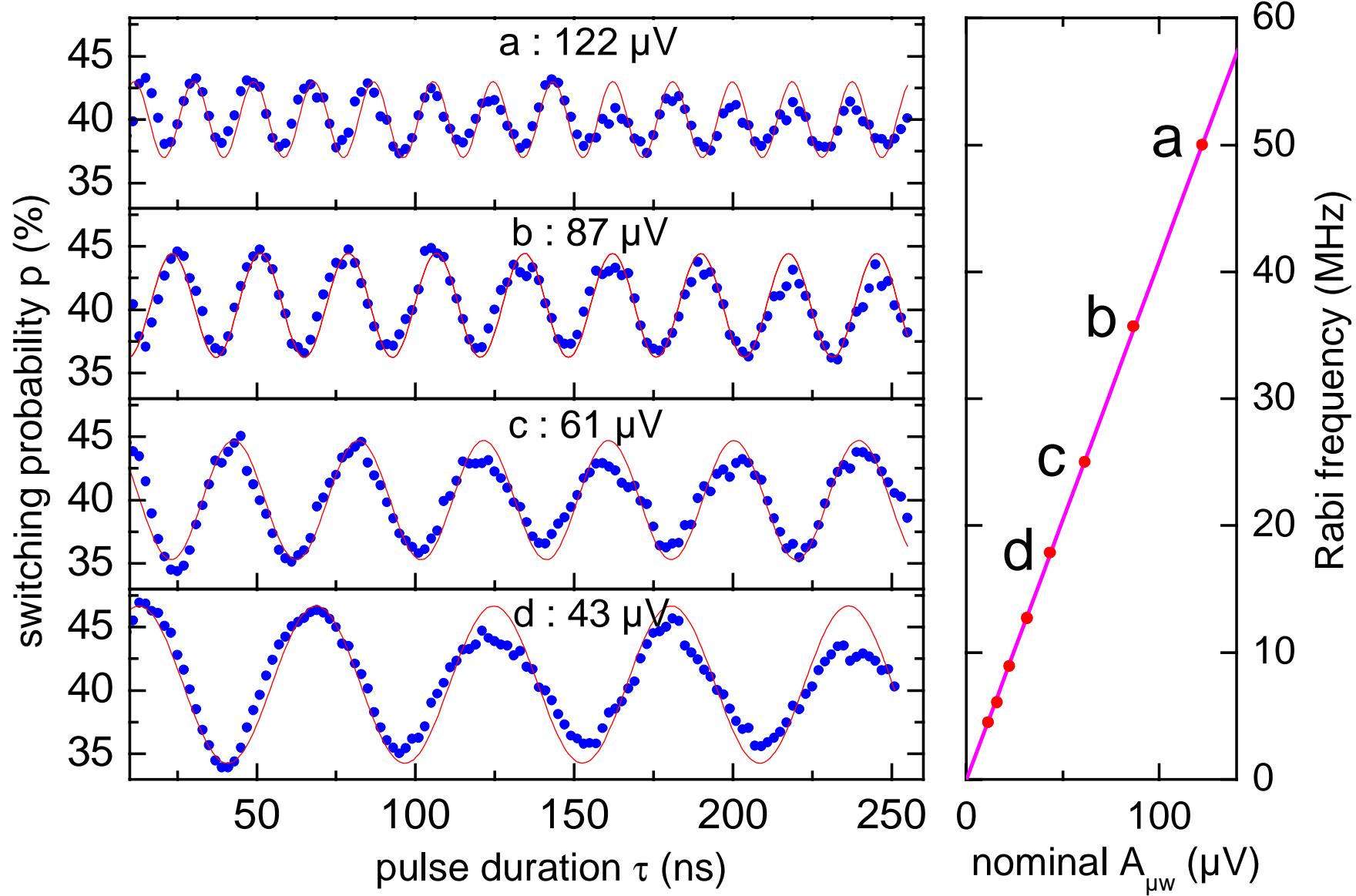
Level spectroscopy $\nu_{01}(N_g, \phi/2\pi)$



1 pulse: quantum state manipulation

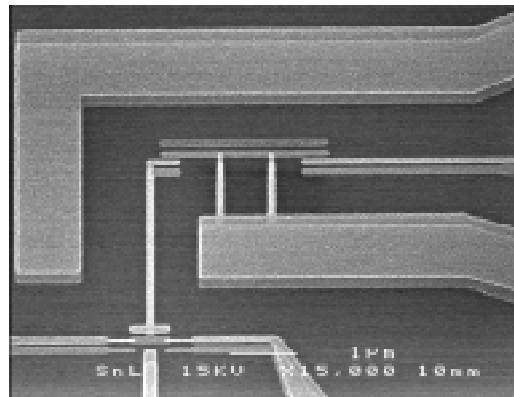


μ w amplitude dependence of Rabi frequency



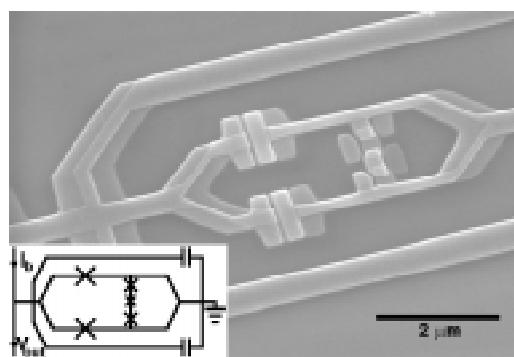
charge qubit

Chalmers U.



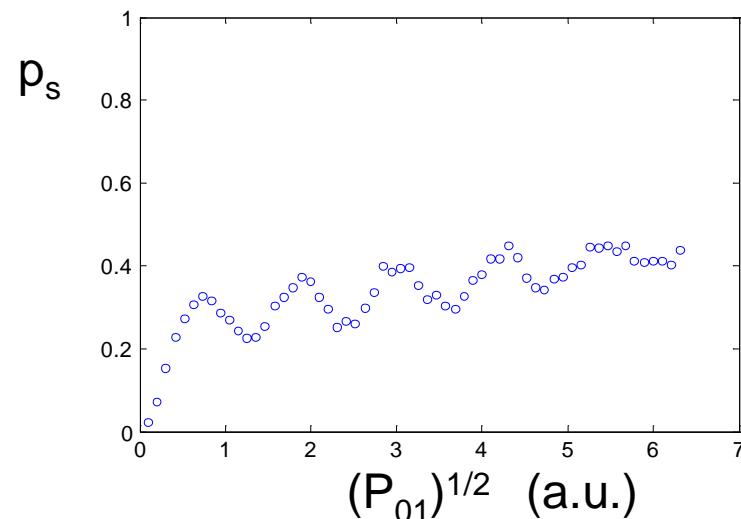
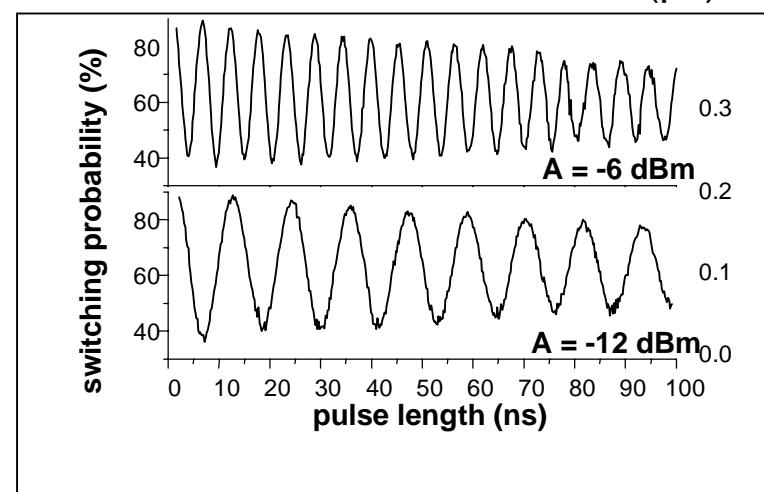
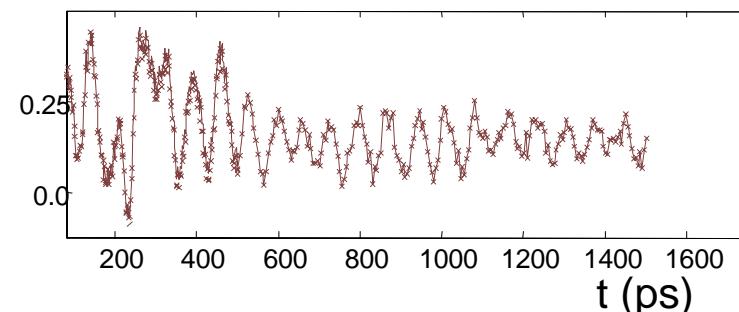
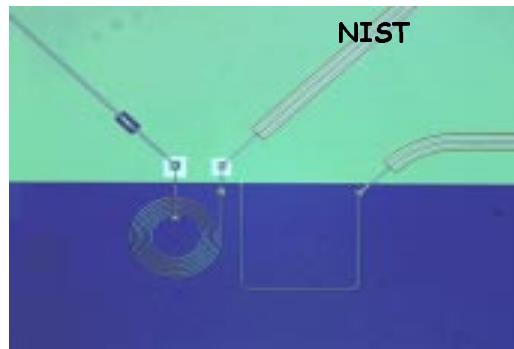
flux qubit

T.U. Delft
(see **hot topic**
K. Harmans)

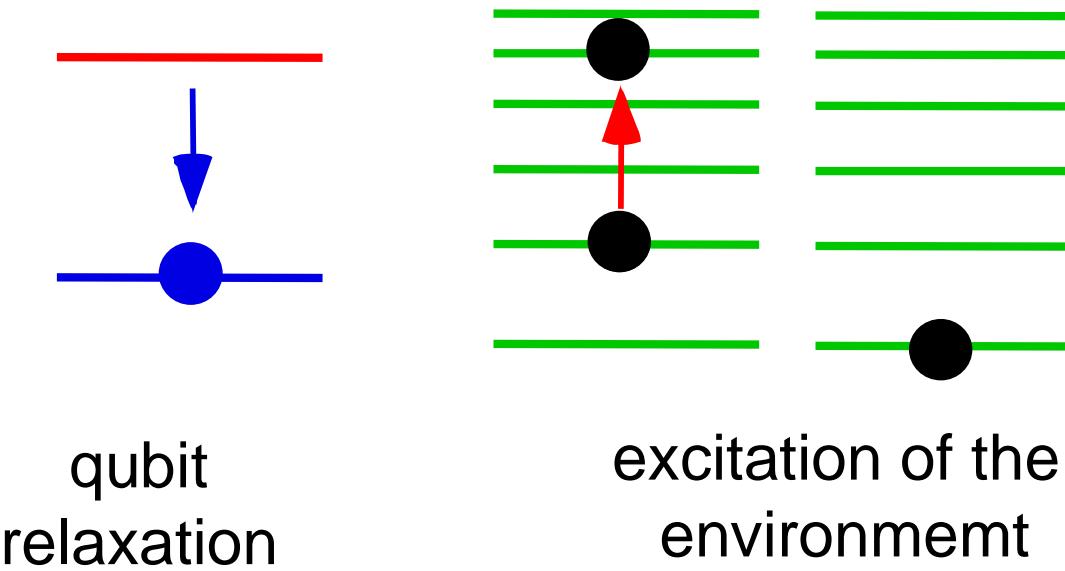


phase qubit

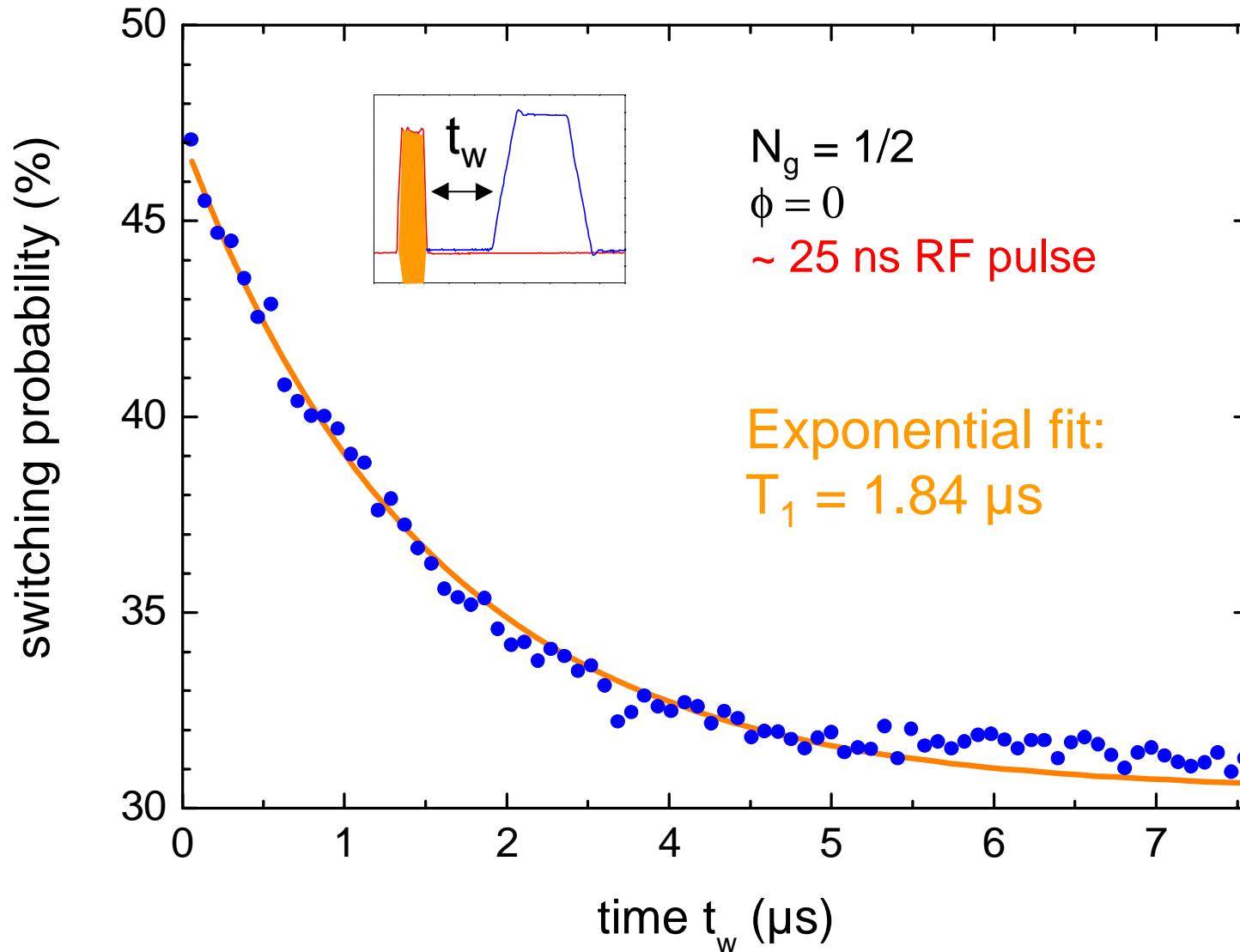
NIST
Martinis et al.



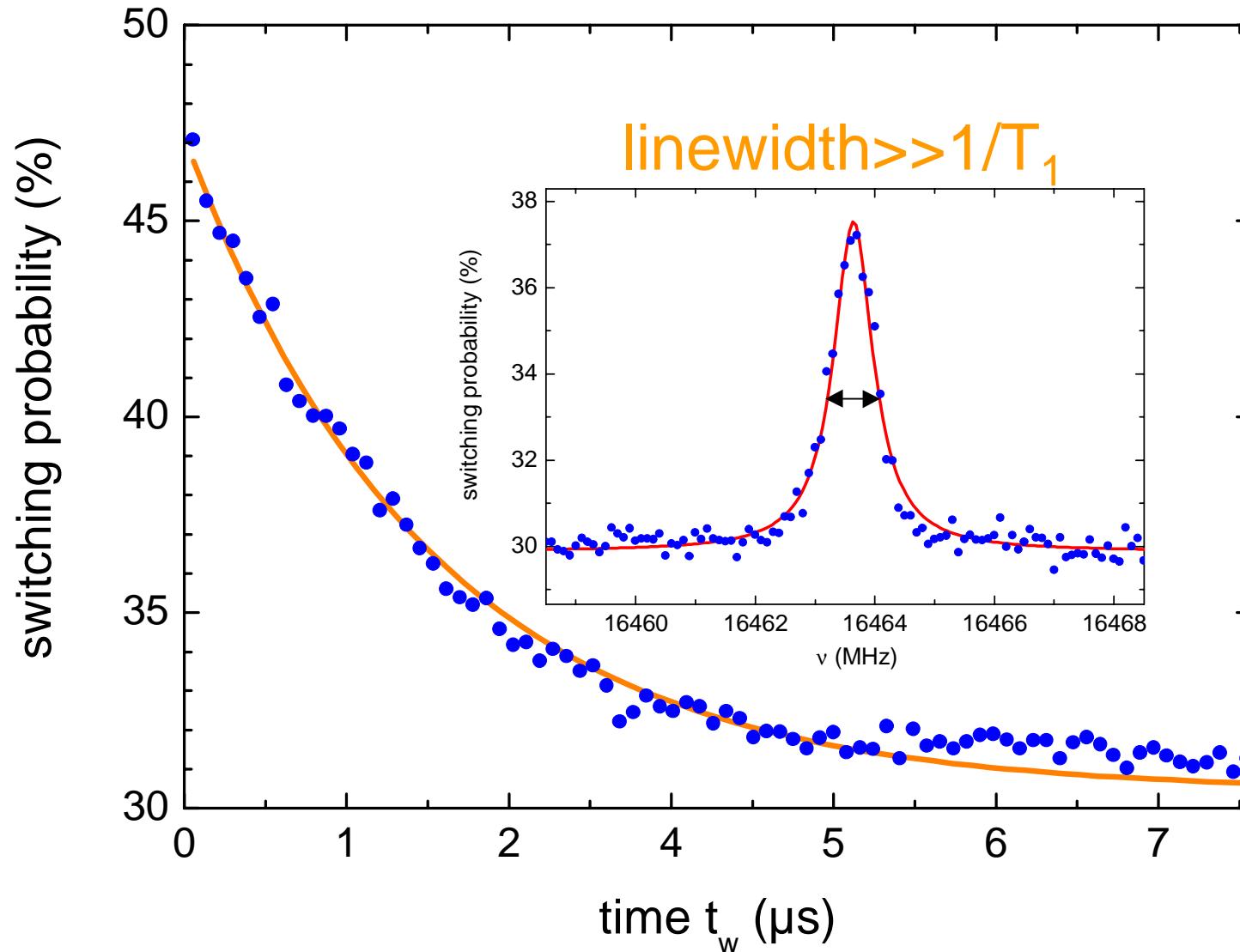
Measurement of the relaxation time



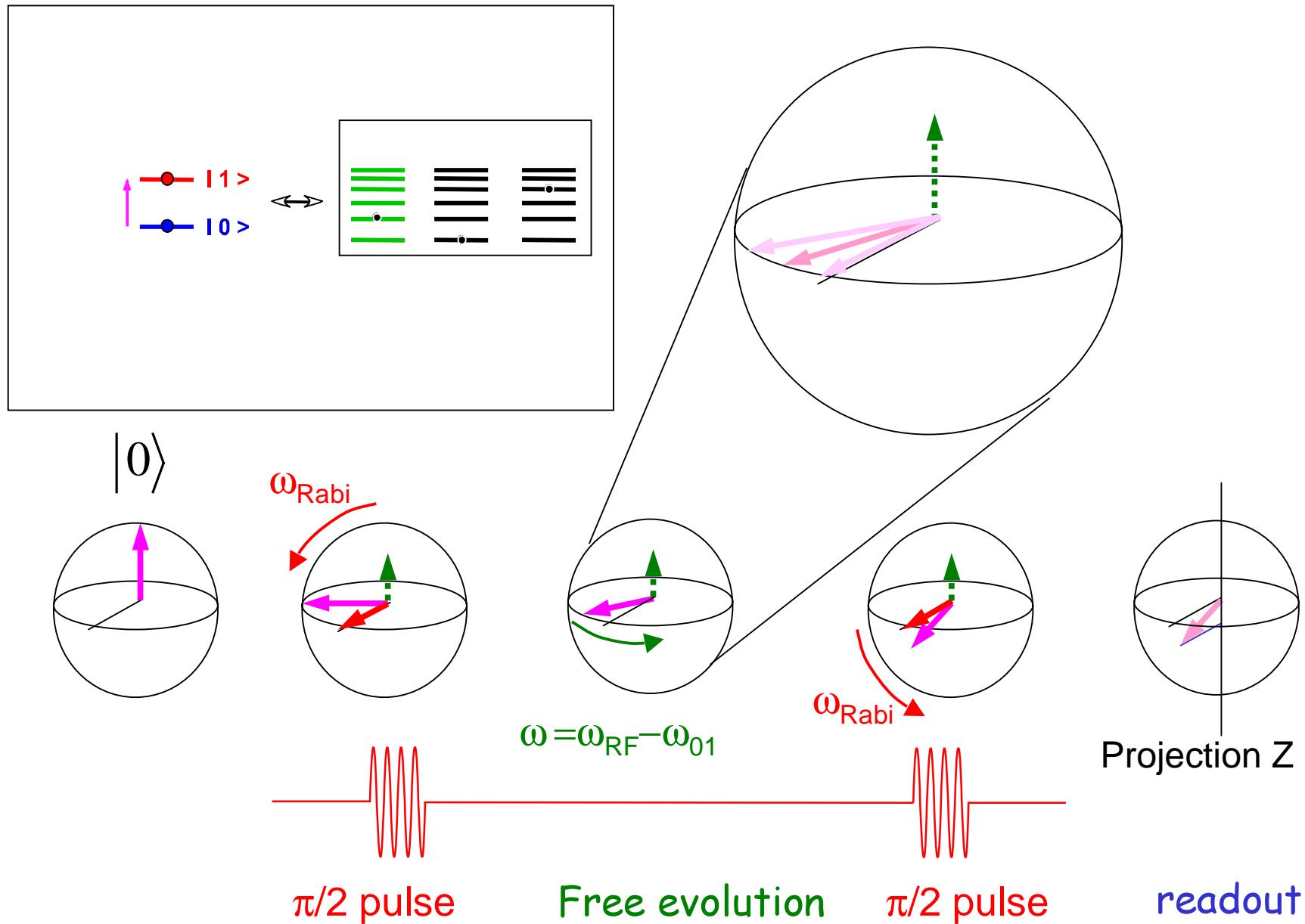
Measurement of the relaxation time



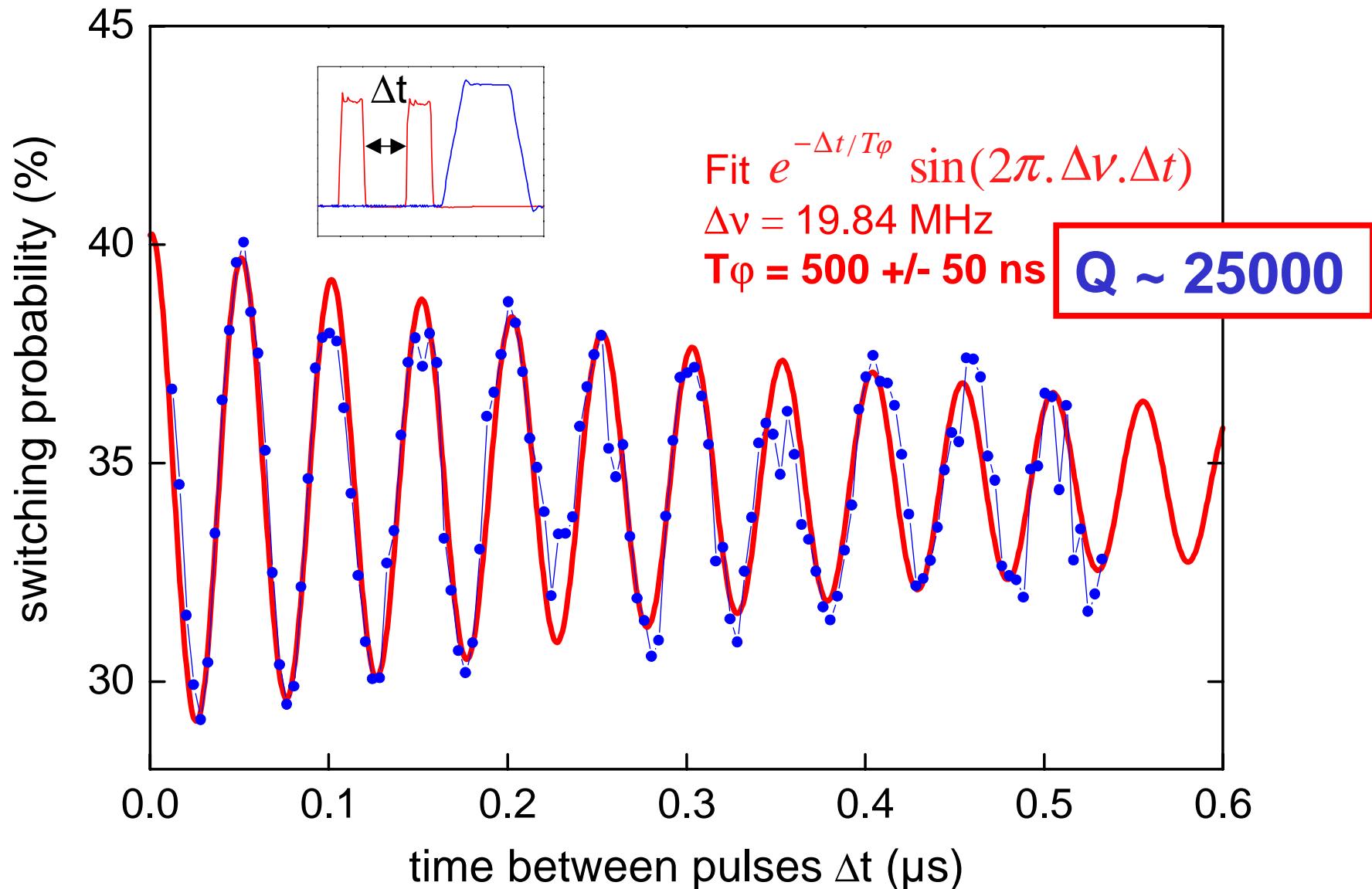
Measurement of the relaxation time



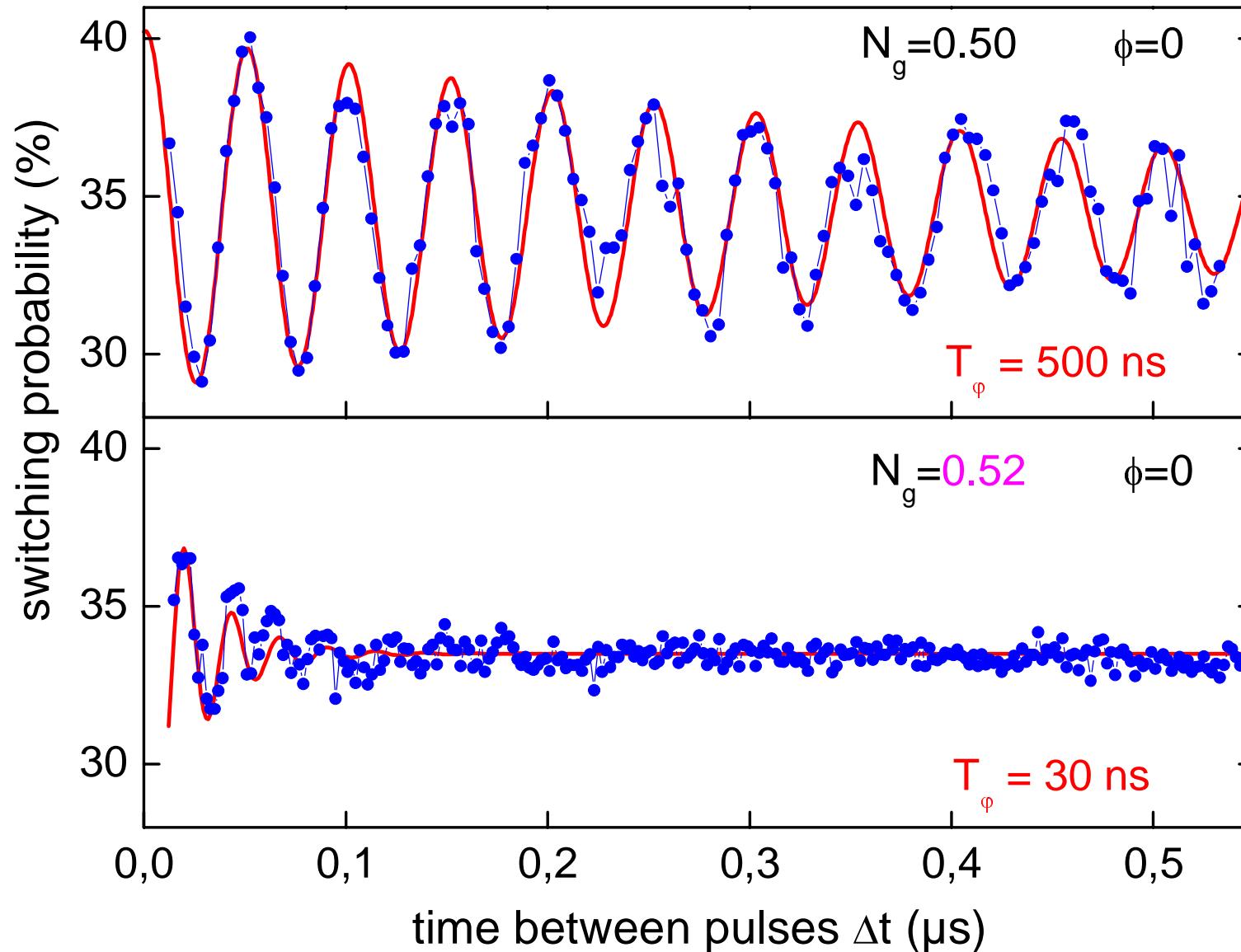
2 pulses: Ramsey interferences



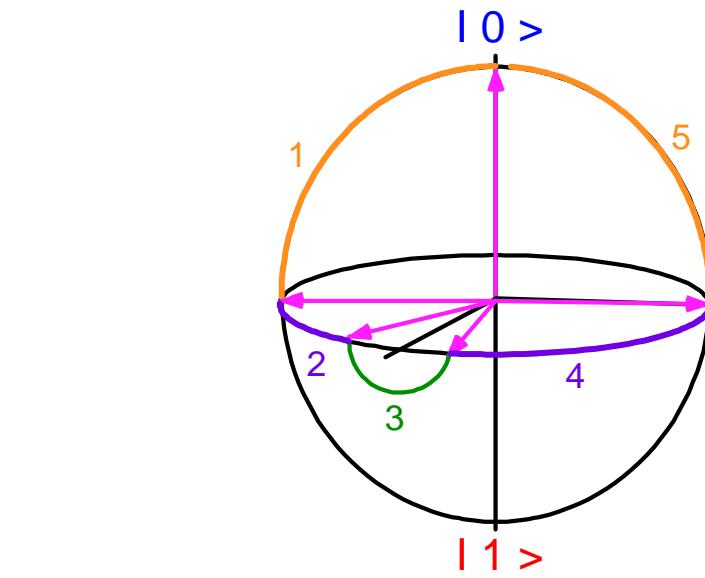
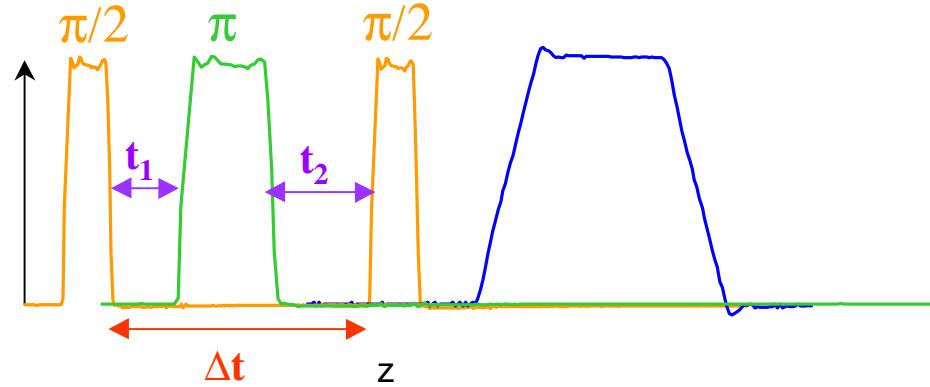
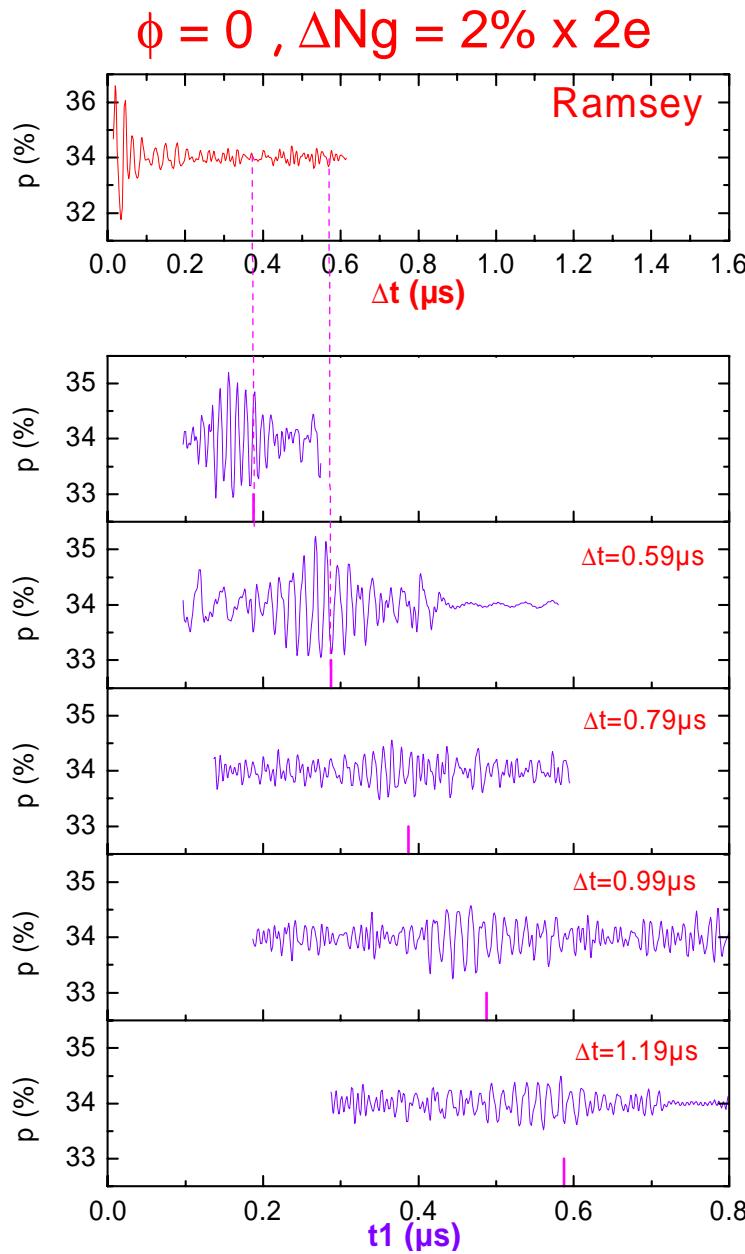
Measurement of the coherence time



Coherence time at the optimal point...and 2% x 2e away

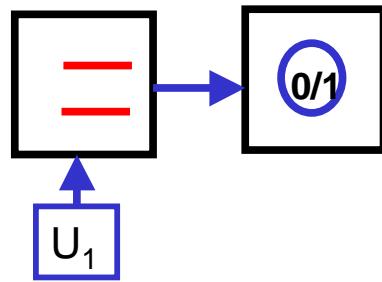


Three pulses: spin-echoes

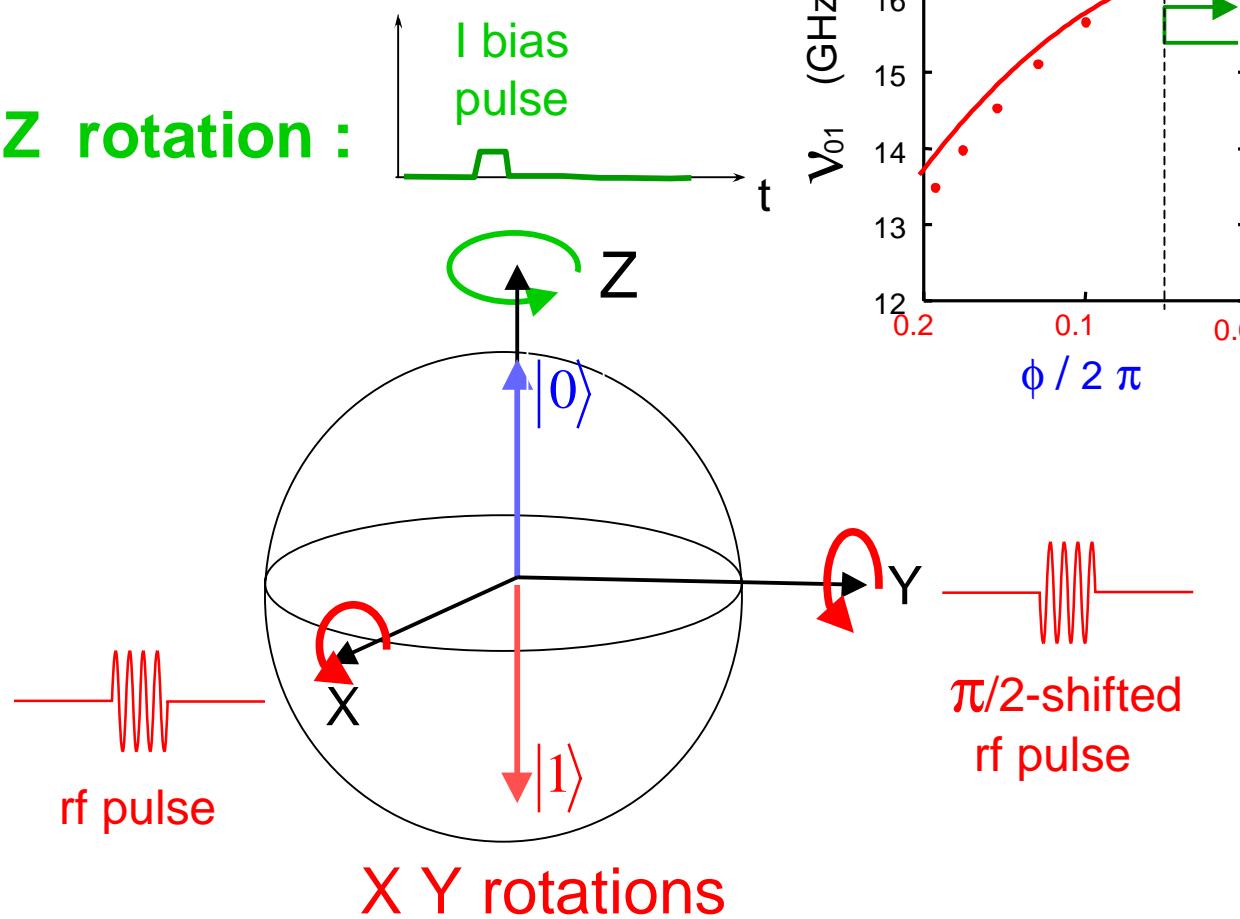


adding ... and removing dephasing

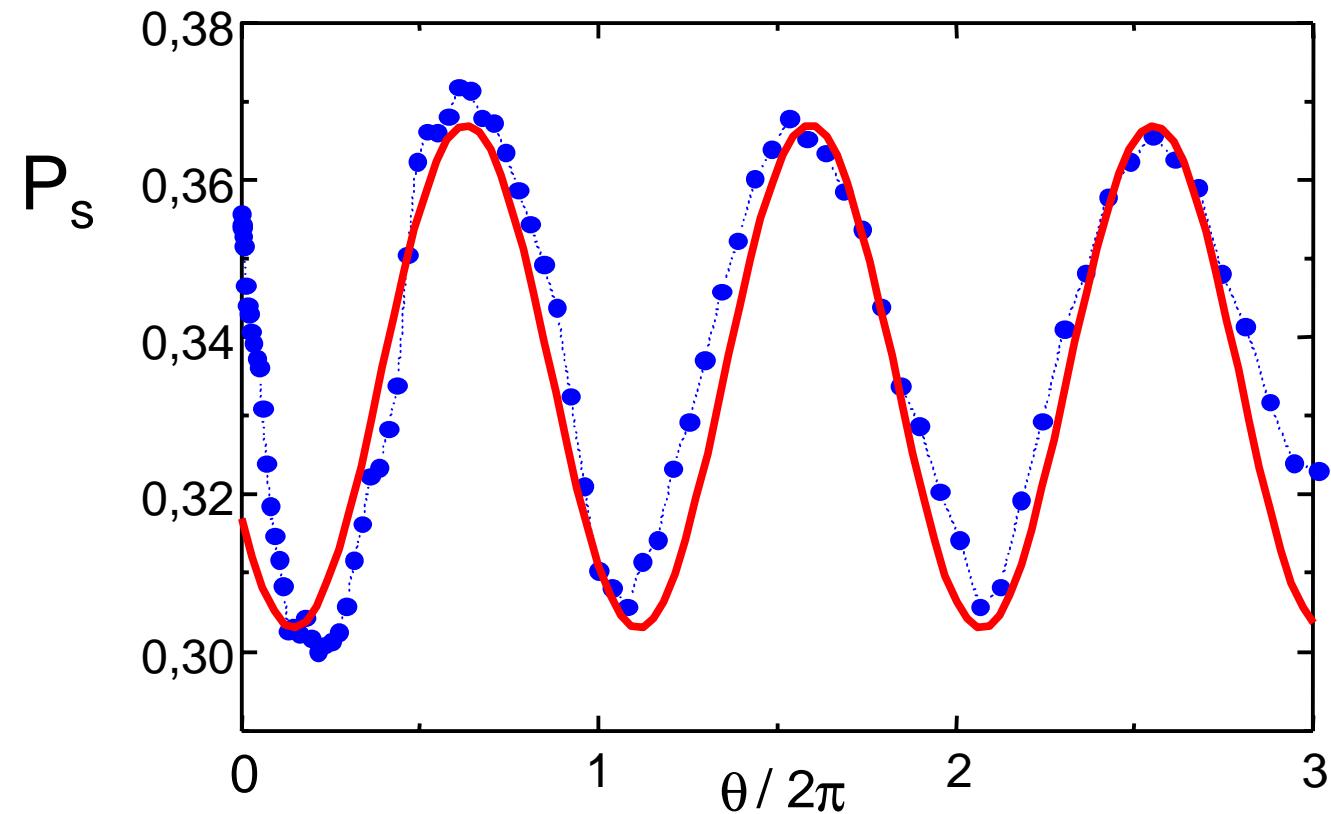
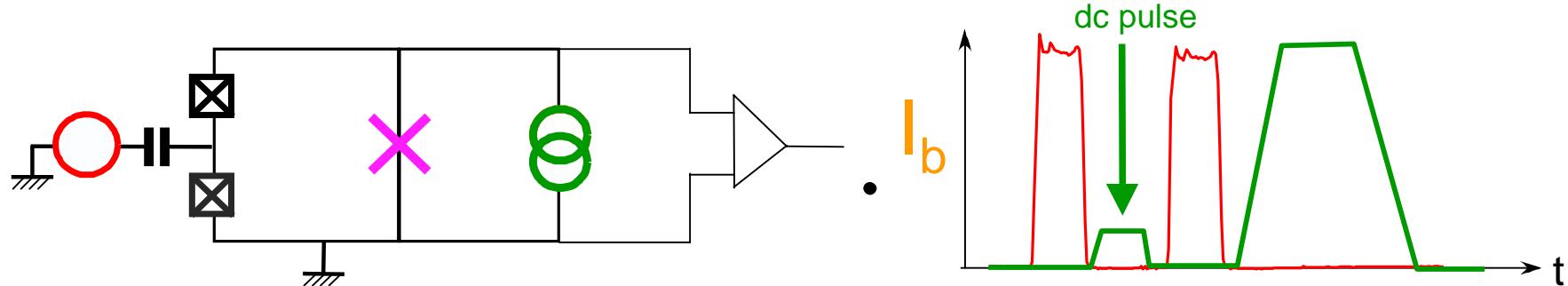
1qubit : full manipulation



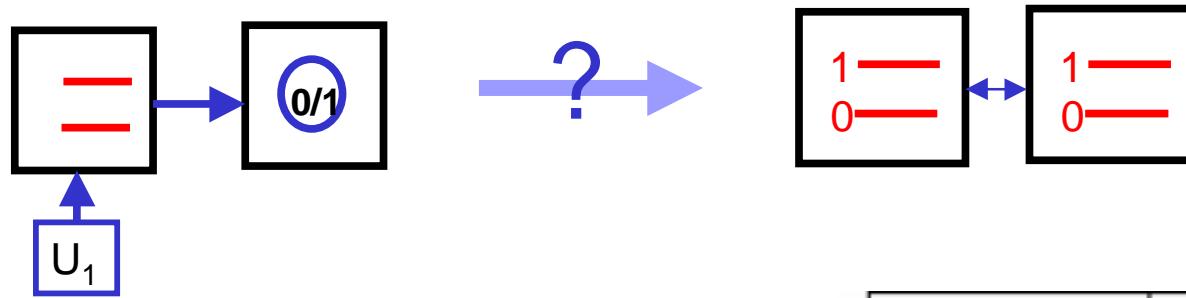
Z rotation :



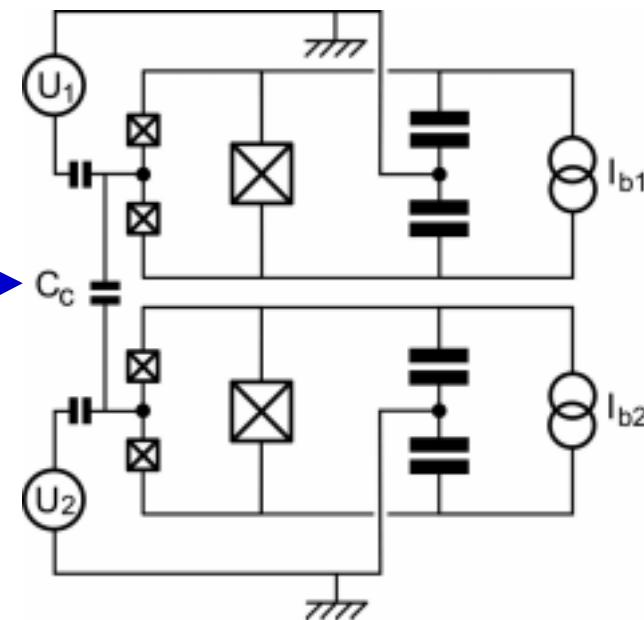
controlled phase-shift



1qubit → 2 qubit gates

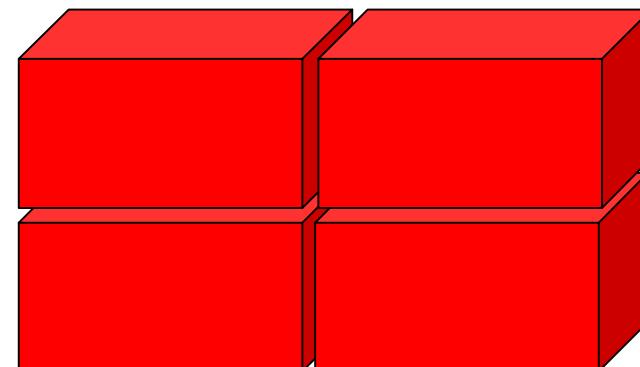
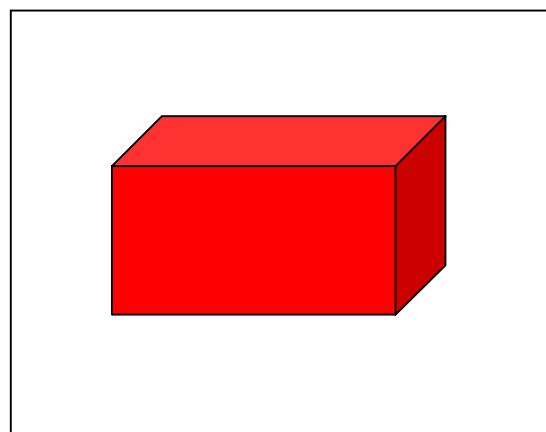
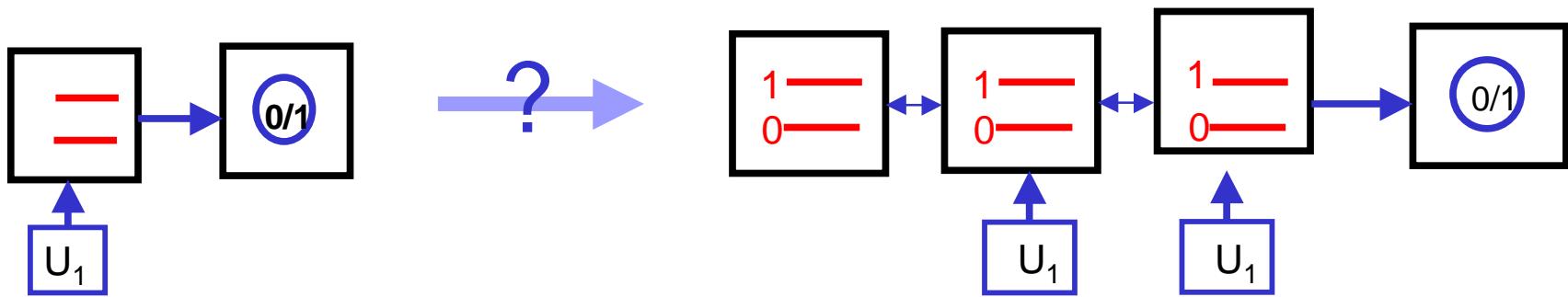


interaction
 $|01\rangle \leftrightarrow |10\rangle$



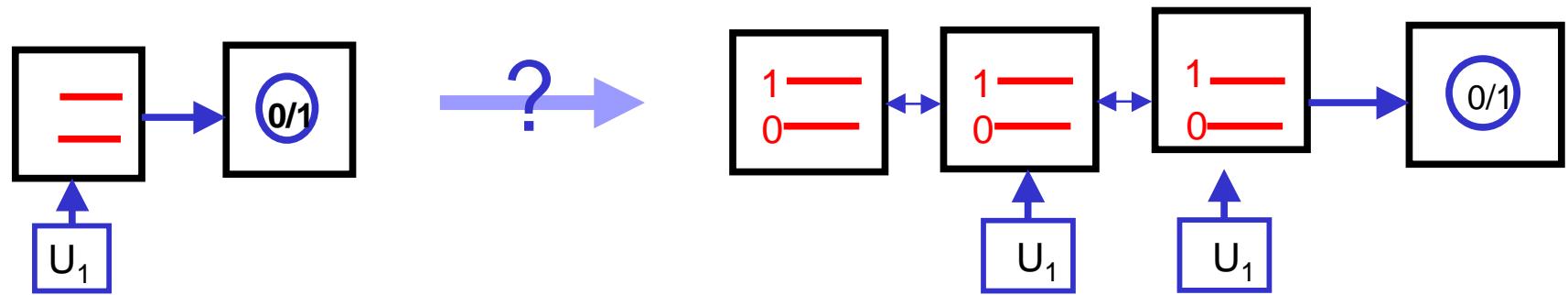
$\sqrt{\text{swap}}$, swap

1qubit → 2 qubit gates → processor



$Q=25000$

1qubit → 2 qubit gates → processor



NEEDED :

- quantum gates
- high fidelity readout(s)
- x100 coherence time



TRY...

QUANTUM ELECTRONICS GROUP

SPEC CEA-Saclay

D. VION
A. COTTET
A. AASSIME
P. JOYEZ
H. POTHIER
M. DEVORET
(now at Yale)

C. URBINA
D. ESTEVE
P. ORFILA
(technician)

and before:
P. LAFARGE
V. BOUCHIAT

